

# FREQUENCY ANALYSIS AND THE TROJAN ASTEROIDS PROBLEM

Frederic Gabern  
Matemàtica Aplicada i Anàlisi  
Universitat de Barcelona



Sieben Heldenköpfe vor 1800 aus Homerwerk. G. Morghen nach J. H. W. Tischbein. Photo © Maier Förlag - GML  
Menelaus, Paris, Diomedes, Odysseus, Nestor, Achilles and Agamemnon.

Recent Trends in Nonlinear Science

Castelló, 24–28 January 2005

## Acknowledgements

- RTNS–2005 organisers.
- J.M. Mondelo, Universitat Autònoma de Barcelona.
- À. Jorba, Universitat de Barcelona.
- P. Robutel, Observatoire de Paris.

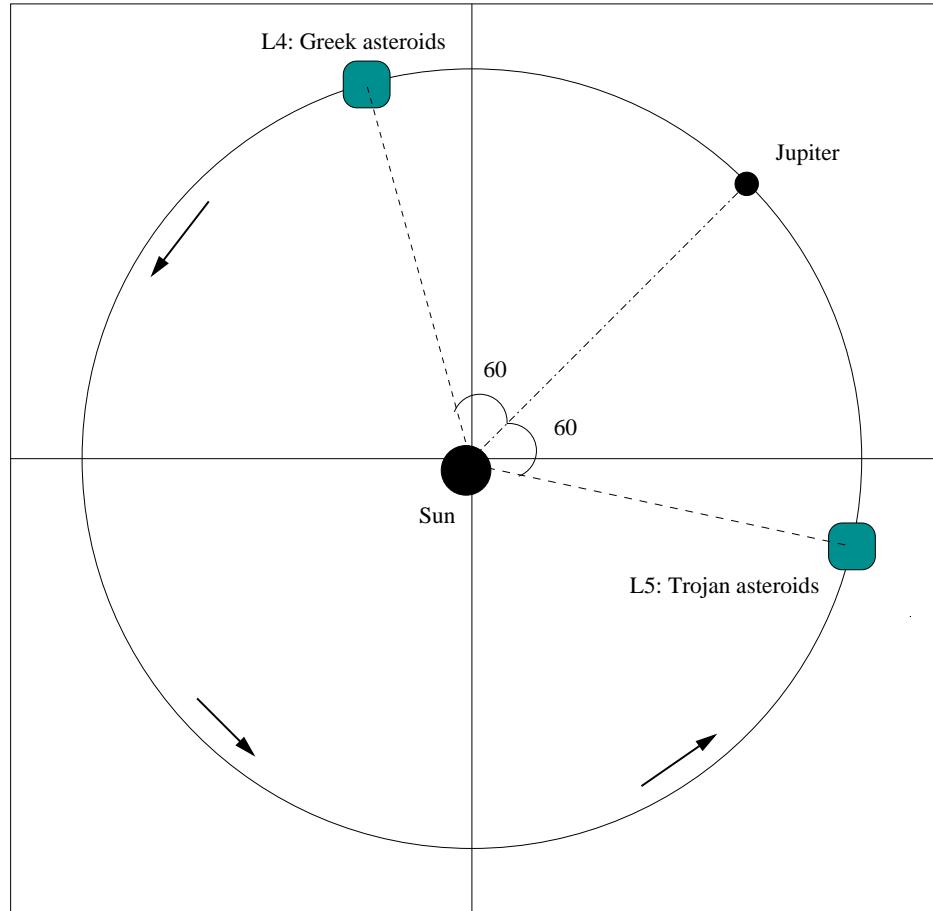
## References

- J. Laskar, “Introduction to frequency map analysis”, In C. Simó, editor, *Hamiltonian Systems with Three or More Degrees of Freedom*, NATO ASI, 134-150. Kluwer, Dordrecht, 1999.
- P. Robutel, and J. Laskar, “Frequency map and global dynamics in the Solar System I”, *Icarus*, 152:4–28, 2001.
- F.G., A. Jorba and P. Robutel, “On the accuracy of Restricted Three-Body Models for the Trojan motion”, *Discrete Contin. Dynam. Systems*, 11(4):843–854, 2004.
- P. Robutel, F.G. and A. Jorba, “The resonant structure of the Jupiter’s Trojan asteroids and its evolution”, *In preparation*, 2005.
- P. Robutel, F.G. and A. Jorba, “The observed Trojans and the global dynamics around the Lagrangian points of the Sun–Jupiter system”, *Celestial Mech.*, to appear. 2005.
- G. Gómez, J.M. Mondelo and C. Simó, “Refined Fourier analysis: procedures, error estimates and applications”, *Preprint*, 2002.
- A. Jorba and M. Zou, “A software package for the numerical integration of ODE by means of high-order Taylor methods”, *Experimental Mathematics*, to appear. 2005.

## Overview

- Introduction.
  - What are the Trojan asteroids?
  - Restricted Three Body Problem.
- Frequency Analysis.
  - Close-to-integrable systems and KAM tori.
  - Numerical approximation of the fundamental frequencies.
- Global dynamics of the Trojans phase space.
  - The Sun–Jupiter–Saturn model.
- Computer Lab:
  - RTBP equations.
  - Numerical integrator of ODE: Taylor method.
  - FURIANE: Frequency Analysis program.
  - Plotting the results: PGPLOT.

## What are the Trojan asteroids? (I)



$L_4$ :

Achilles



$L_5$ :

Patroclus



Hector



Aeneas



Nestor



Memnon



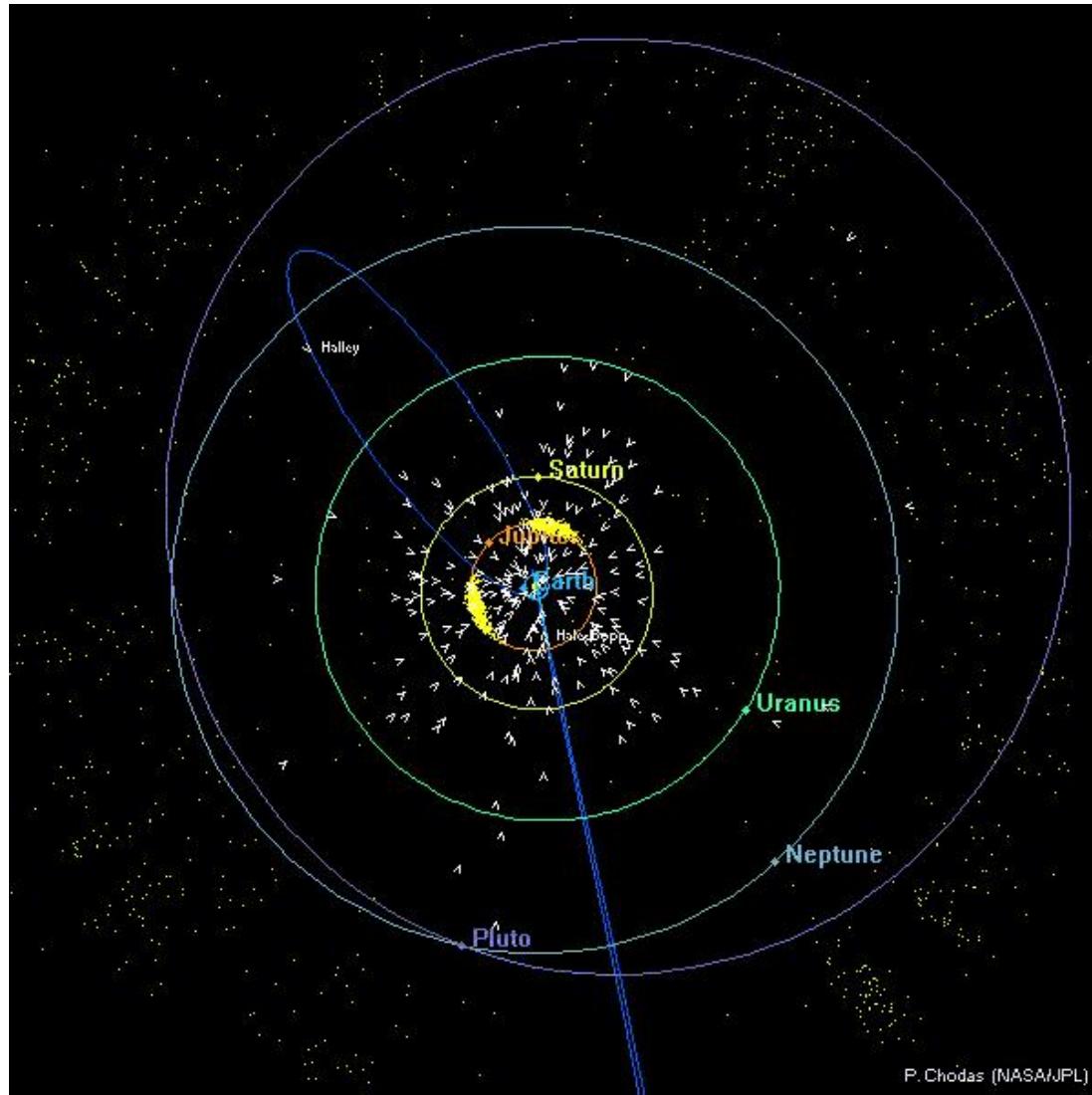
Agamemnon



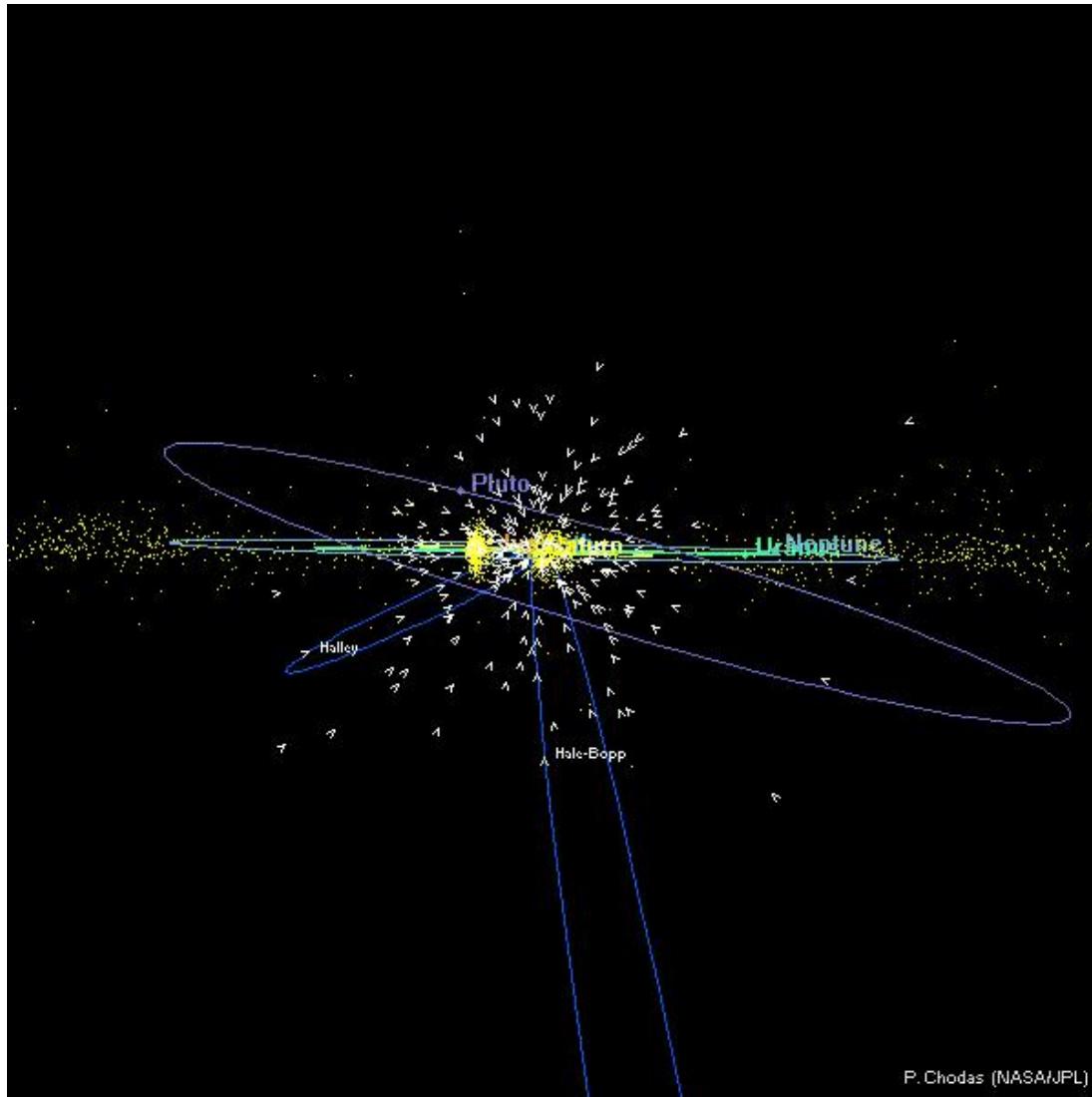
Paris



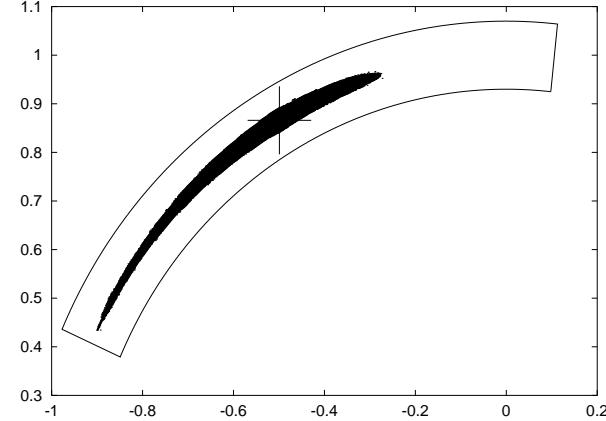
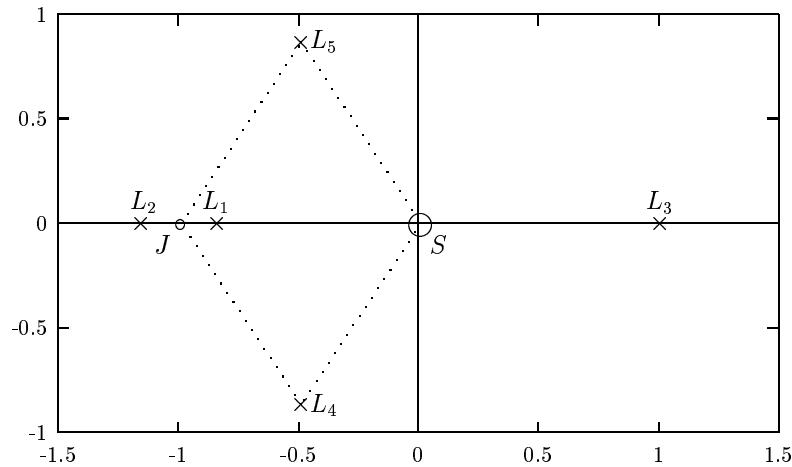
## What are the Trojan asteroids? (II)



## What are the Trojan asteroids? (III)

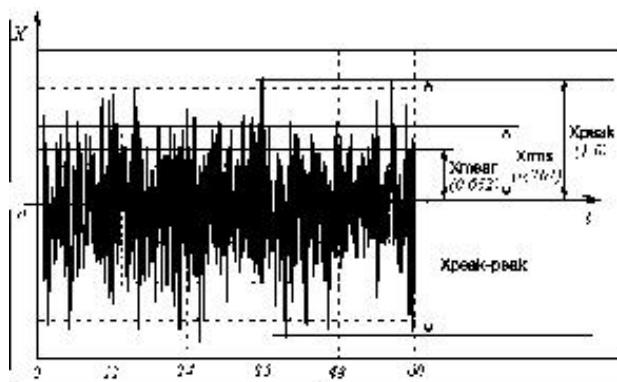


## The Restricted Three Body Problem



$$\begin{aligned} H = & \frac{1}{2} (p_x^2 + p_y^2 + p_z^2) + (yp_x - xp_y) \\ & - \frac{1 - \mu}{\sqrt{(x - \mu)^2 + y^2 + z^2}} - \frac{\mu}{\sqrt{(x - \mu + 1)^2 + y^2 + z^2}} \end{aligned}$$

## FREQUENCY ANALYSIS



## Frequency Analysis: Close-to-integrable systems and KAM tori (I)

$$H(I, \theta) = H_0(I) + \varepsilon H_1(I, \theta), \quad (I, \theta) \in B^n \times \mathbb{T}^n, \quad B^n \subset \mathbb{R}^n$$

- For  $\varepsilon = 0$ :

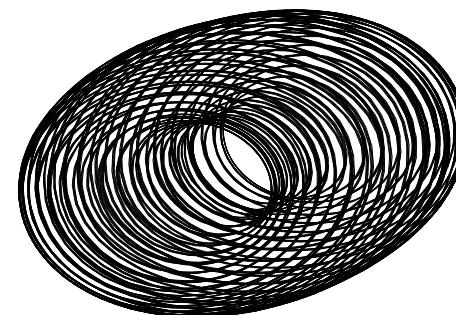
$$\dot{I}_j = 0, \quad \dot{\theta}_j = \frac{\partial H_0}{\partial I_j} \equiv \omega_j(I); \quad j = 1, \dots, n$$

If  $\det \left( \frac{\partial \omega(I)}{\partial I} \right) \neq 0$  (non-degenerate system), the *frequency* map

$$\begin{aligned} F : \quad B^n &\longrightarrow \mathbb{R}^n \\ I &\longmapsto \omega(I) \end{aligned}$$

is a diffeomorphism on its image.

Motion in phase space takes place on *invariant tori*.



## Frequency Analysis: Close-to-integrable systems and KAM tori (II)

$$H(I, \theta) = H_0(I) + \varepsilon H_1(I, \theta), \quad (I, \theta) \in B^n \times \mathbb{T}^n, \quad B^n \subset \mathbb{R}^n$$

- For  $\varepsilon \neq 0$ : The *KAM theorem* asserts that, for sufficiently small  $\varepsilon$ , there exists a Cantor set  $\Omega_\varepsilon \ni \omega$ , for which the invariant tori of the unperturbed system persist: *KAM tori*.

For  $\omega \in \Omega_\varepsilon$ , the solution lies on a torus and can be parametrized, for instance, by a Fourier series:

$$z^j(t) = z_0^j e^{i\omega_j t} + \sum_k a_k^j(\omega) e^{i\langle k, \omega \rangle t}, \quad j = 1, \dots, n$$

If we fix  $\theta \in \mathbb{T}^n$  to some value  $\theta = \theta_0$ , we obtain the frequency map on  $B^n$ :

$$F_{\theta_0} : B^n \longrightarrow \mathbb{R}^n$$

## Frequency Analysis: Quasi-periodic approximation (I)

Let

$$f(t) = \sum_{k \in \mathbb{Z}^m} a_k e^{i \langle k, \omega \rangle t}, \quad a_k \in \mathbb{C},$$

be a quasi-periodic function for which we know a table of equidistant values in the time span  $[0, T]$ :  $f(t_k)$ ,  $t_k \in [0, T]$ ,  $k = 1, \dots, N$ .

The frequency analysis algorithm provides (numerically) the values of the:

- Frequencies:  $\tilde{\omega}_k$
- Amplitudes:  $\tilde{a}_k$

of a function

$$\tilde{f}(t) = \sum \tilde{a}_k e^{i \tilde{\omega}_k t}$$

that approximates  $f(t)$  in  $[0, T]$ .

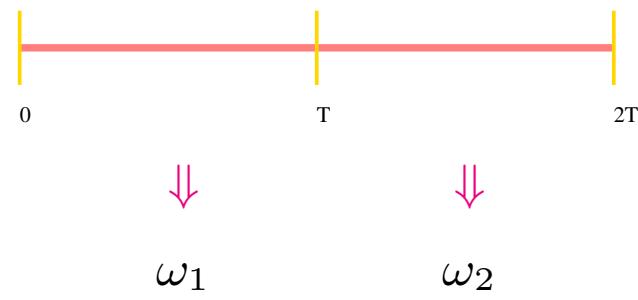
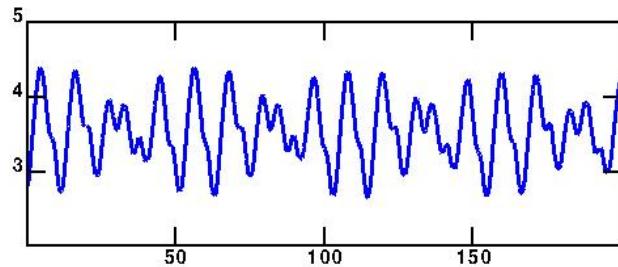
## Frequency Analysis: Quasi-periodic approximation (II)

Phase space of a close-to-integrable system:



Quasi-Periodic Motions

Frequency decomposition of particular trajectories



Diffusion index:

$$\delta = \left| 1 - \frac{\omega_2}{\omega_1} \right|$$

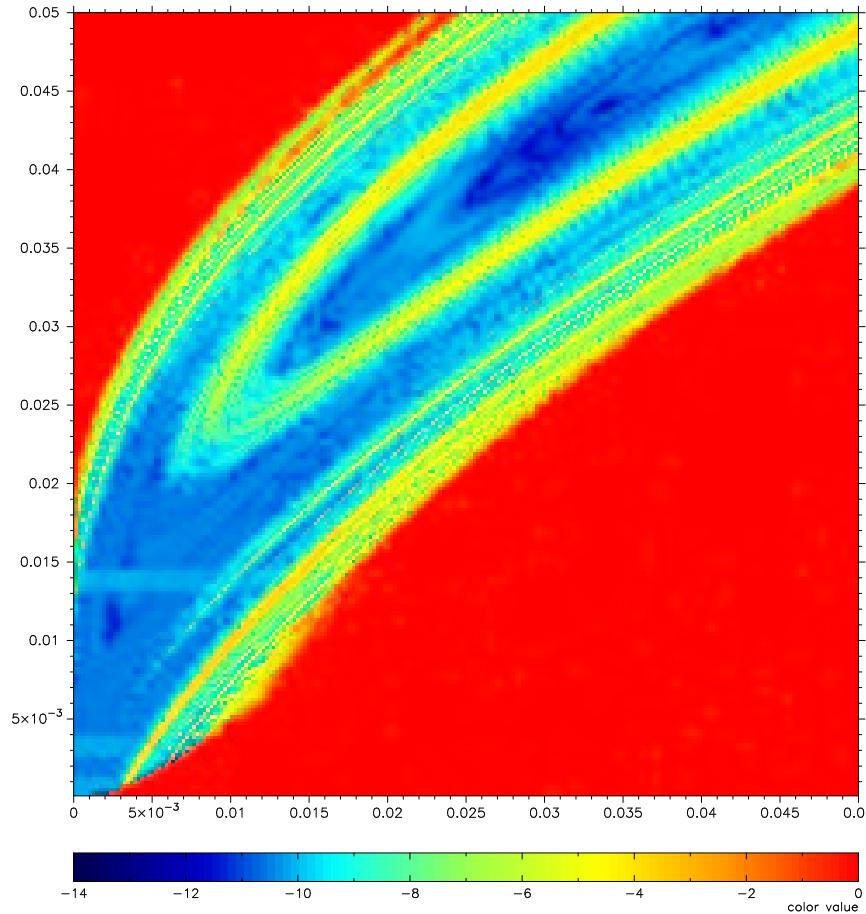
- $\delta$  gives an estimation of the “chaoticity” of the particular orbit.
- If the trajectory associated to an initial condition is quasi-periodic,  $\delta$  is zero.

## Frequency Analysis: Example I

RTBP near  $L_5$ :  $H(I, \varphi) = \nu_1 I_1 + \nu_2 I_2 + \sum_{l=3}^{+\infty} h_l(I, \varphi)$

Diffusion index:

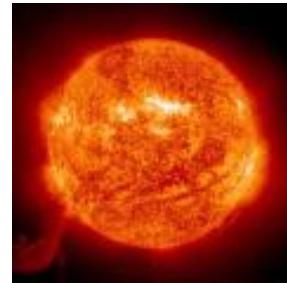
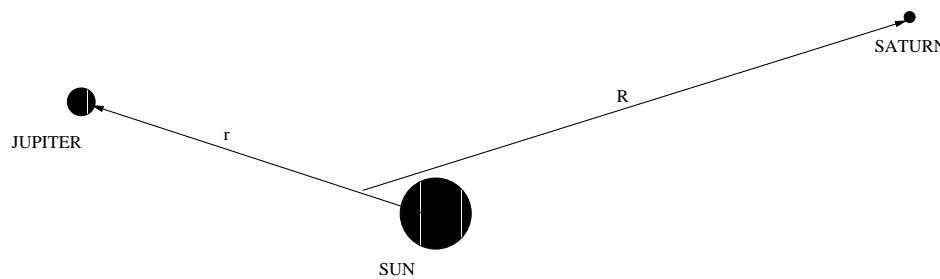
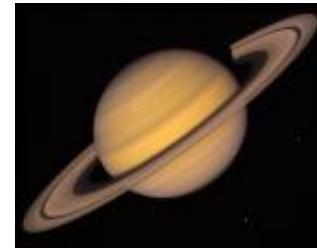
$$\sigma = \log \left| 1 - \frac{\omega_2}{\omega_1} \right|$$



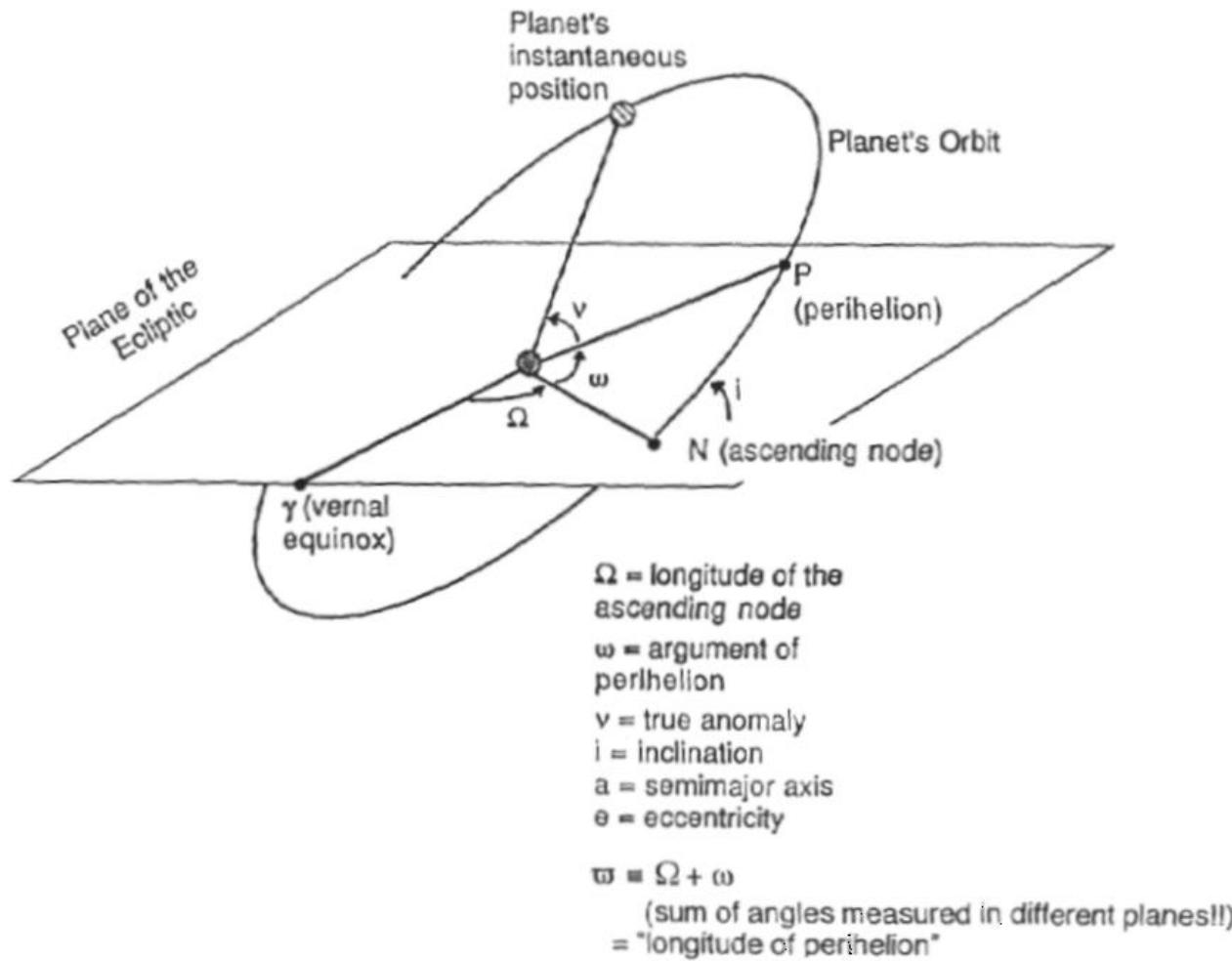
**GLOBAL DYNAMICS  
OF THE  
TROJANS PHASE SPACE**



# Trojan Global dynamics: Sun–Jupiter–Saturn Model



## Trojan Global dynamics: Orbital Elements



## Trojan Global dynamics: Basic Frequencies of the planets (I)

Time window: 10 Million years.

Motion of the planets can be considered quasi-periodic.

Apply frequency analysis to:

$$\begin{aligned}\alpha_p(t) &= a_p(t) \exp(i\lambda_p(t)), \\ \beta_p(t) &= e_p(t) \exp(i\varpi_p(t)), \\ \gamma_p(t) &= \sin\left(\frac{I_p(t)}{2}\right) \exp(i\Omega_p(t)),\end{aligned}$$

where

- $p = 5 \rightarrow$  Jupiter       $p = 6 \rightarrow$  Saturn.
- $\lambda = \varpi + M$ : mean longitude ( $M$ : mean anomaly).
- $\varpi = \Omega + \omega$ : longitude of the perihelion.

## Trojan Global dynamics: Basic Frequencies of the planets (II)

	frequency ("/yr)	Period (yr)
$n_5$	109254.63165	11.8622
$n_6$	43995.34975	29.4577
$g_5$	4.02760	321780
$g_6$	28.00657	46274.9
$s_6$	-26.03912	49771.3

- $\{n_5, n_6\}$ : Proper Mean Motion of the planets.
- $\{g_5, g_6, s_6\}$ : Secular frequencies.

## Trojan Global dynamics: Frequencies of the asteroids (I)

Time window: 5 Million years.

Apply frequency analysis to:

$$\begin{aligned}\alpha_j(t) &= a_j(t) \exp(i(\lambda_j(t) - \lambda_5(t))), \\ \beta_j(t) &= e_j(t) \exp(i\varpi_j(t)), \\ \gamma_j(t) &= \sin\left(\frac{I_j(t)}{2}\right) \exp(i\Omega_j(t)).\end{aligned}$$

FREQUENCY MAP:

$$\begin{aligned}F_{\theta_0} &: (a, e, I) \mapsto (\nu, g, s), \\ \theta_0 &= (\lambda_0, \varpi_0, \Omega_0)\end{aligned}$$

## Trojan Global dynamics: Frequencies of the asteroids (II)

- Fixed phase:

$$\begin{aligned}\lambda_0 &= \lambda_5 \mp \frac{\pi}{3}, \quad (L_4, L_5) \\ \varpi_0 &= \varpi_5 \mp \frac{\pi}{3}, \quad (L_4, L_5) \\ \Omega_0 &= \Omega_5.\end{aligned}$$

- Fixed inclination,  $I = I_5$ :

FREQUENCY MAP:

$$F_{(I_5, \theta_0)} : (a, e) \mapsto (\nu, g, s).$$

## Trojan Global dynamics: Diffusion

Frequency of libration:

$$\nu \rightsquigarrow [0, T]$$

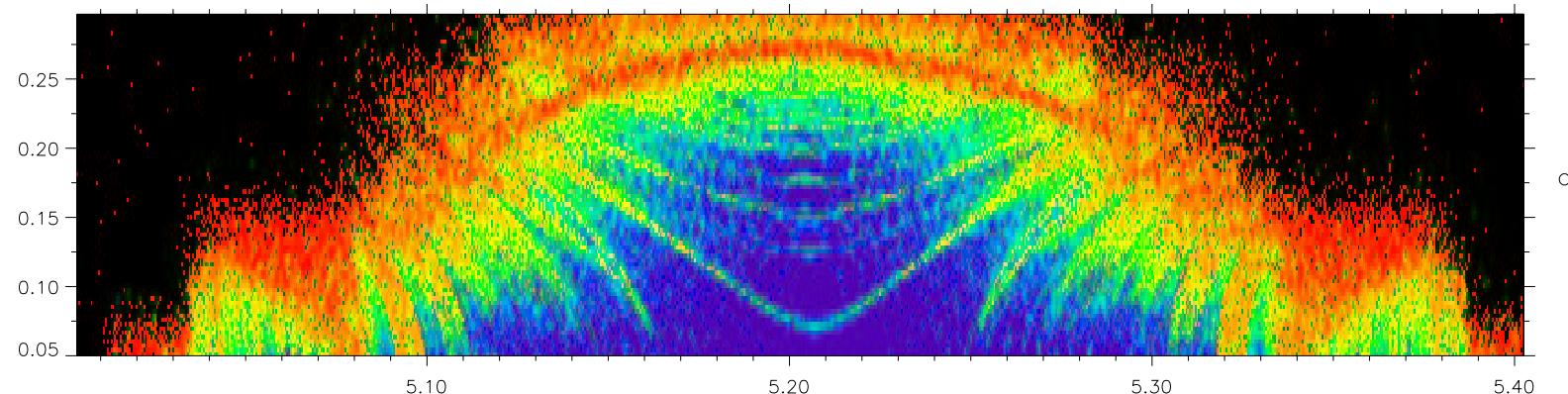
$$\nu' \rightsquigarrow [T, 2T]$$

Diffusion index:

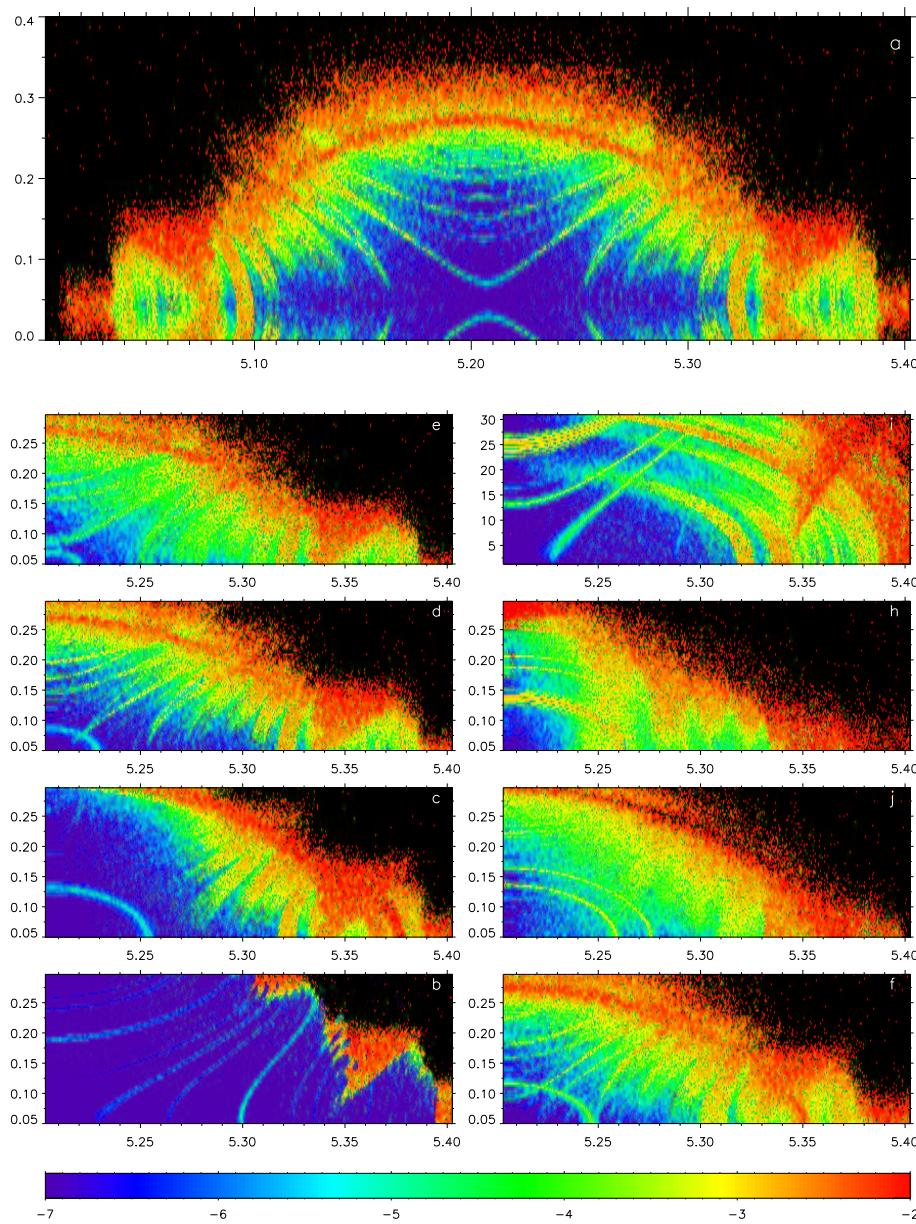
$$\delta\nu = 1 - \frac{\nu'}{\nu}$$

Color code:

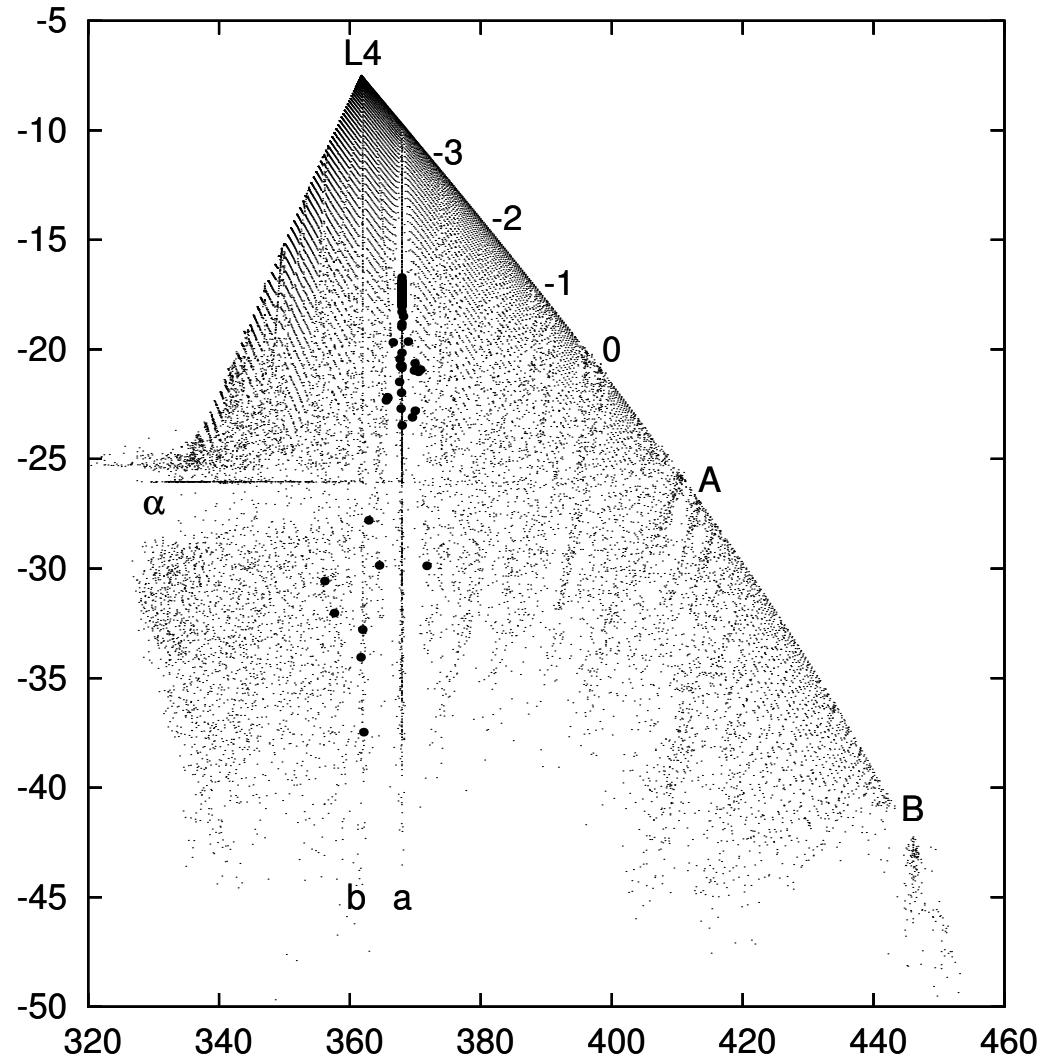
- $\delta\nu > 10^{-2}$  : Strongly irregular motion. [RED]
- $\delta\nu < 10^{-7}$  : Close-to quasi-periodic motion. [BLUE]
- [BLACK] : Escaping orbits.



D  
Y  
N  
A  
M  
I  
C  
S



## Trojan Global dynamics: Frequency Space ( $g, s$ )



## Global dynamics: Observed Trojans (I)

How to locate the observed Trojans in the global dynamical pictures?

- FREQUENCY MAP:  $F_{\theta_0} : (a, e, I) \mapsto (\nu, g, s)$ .
- PHASE PROBLEM: For every observed Trojan  $\theta_j \neq \theta_0$ .
- IDEA: Look at the *Frequency Space*.

For the Trojan asteroid  $j$ :

$$(a_j, e_j, I_j, \lambda_j, \varpi_j, \Omega_j) \implies F_{\theta_j}(a_j, e_j, I_j) = (\nu_j, g_j, s_j)$$

Let  $\mathcal{D}$  be the domain of initial conditions:  $(a, e, I) \in \mathcal{D}$  and  $\mathcal{F} = F_{\theta_0}(\mathcal{D})$ .

OBS:  $F_{\theta_j}(a_j, e_j, I_j) \in \mathcal{F}$ . Then, we define (well defined if  $H$  integrable)

$$(\tilde{a}_j, \tilde{e}_j, \tilde{I}_j) = F_{\theta_0}^{-1} \circ F_{\theta_j}(a_j, e_j, I_j).$$

2 ≠ trajectories lie on the SAME TORUS  $\Rightarrow$  DYNAMICALLY EQUIVALENT

$\hookrightarrow (\tilde{a}_j, \tilde{e}_j, \tilde{I}_j, \theta_0)$  vs.  $(a_j, e_j, I_j, \theta_j)$

FREQUENCY VECTOR  $\iff$  INVARIANT TORUS

## Global dynamics: Observed Trojans (II)

In practice, we find two main difficulties:

1.  $H$  is not integrable in an open subset of the phase space

$\Downarrow$

$F_{\theta_0}^{-1}$  is not invertible.

2. The domain of initial conditions  $\mathcal{D}$  is a discrete set:  $\tilde{\mathcal{D}}$

$\Downarrow$

Discrete Frequency Space:  $\tilde{\mathcal{F}} = F_{\theta_0}(\tilde{\mathcal{D}})$

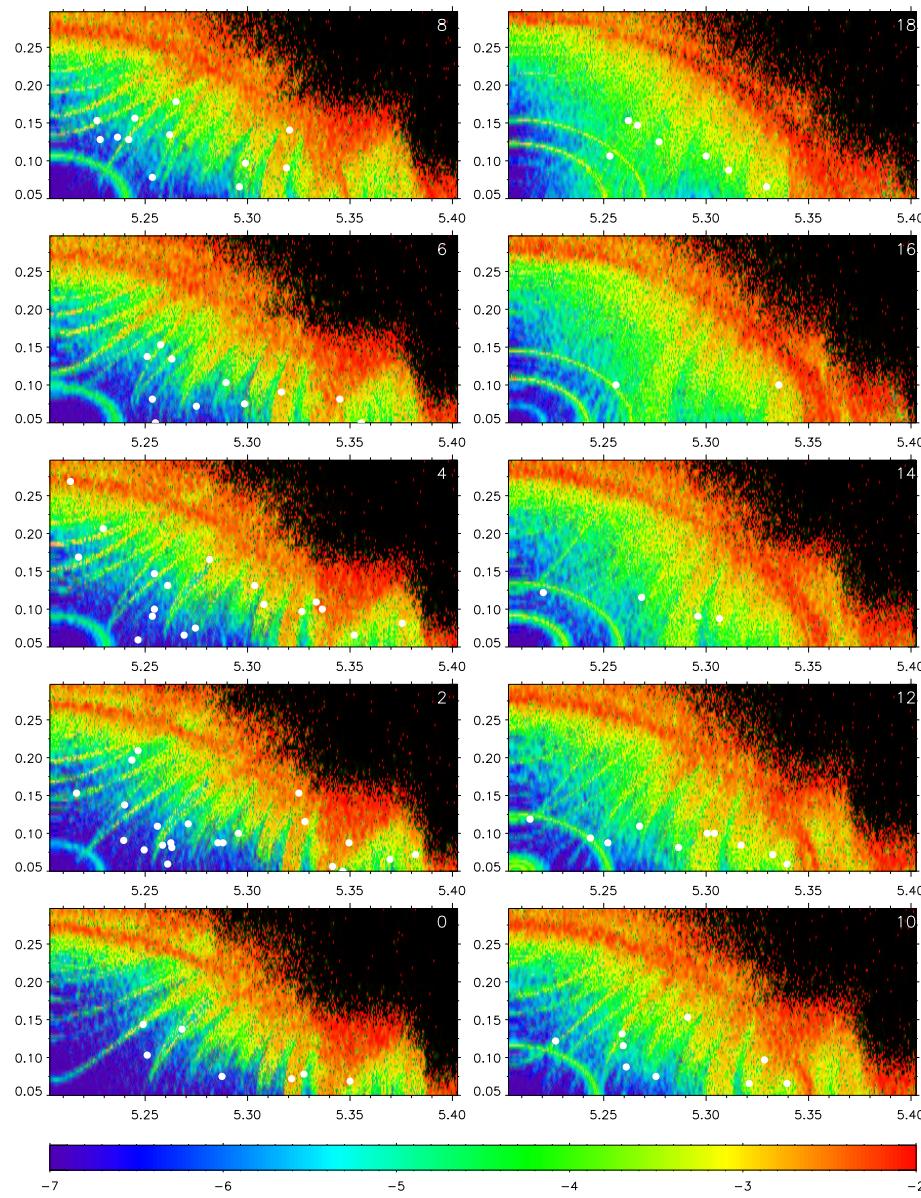
For a given Trojan, we approximate  $(\tilde{a}_j, \tilde{e}_j, \tilde{I}_j) \approx (a_j^*, e_j^*, I^*)$  such that

$$d_{j,0} = \|F_{\theta_j}(a_j, e_j, I_j) - F_{\theta_0}(a^*, e^*, I^*)\|_2$$

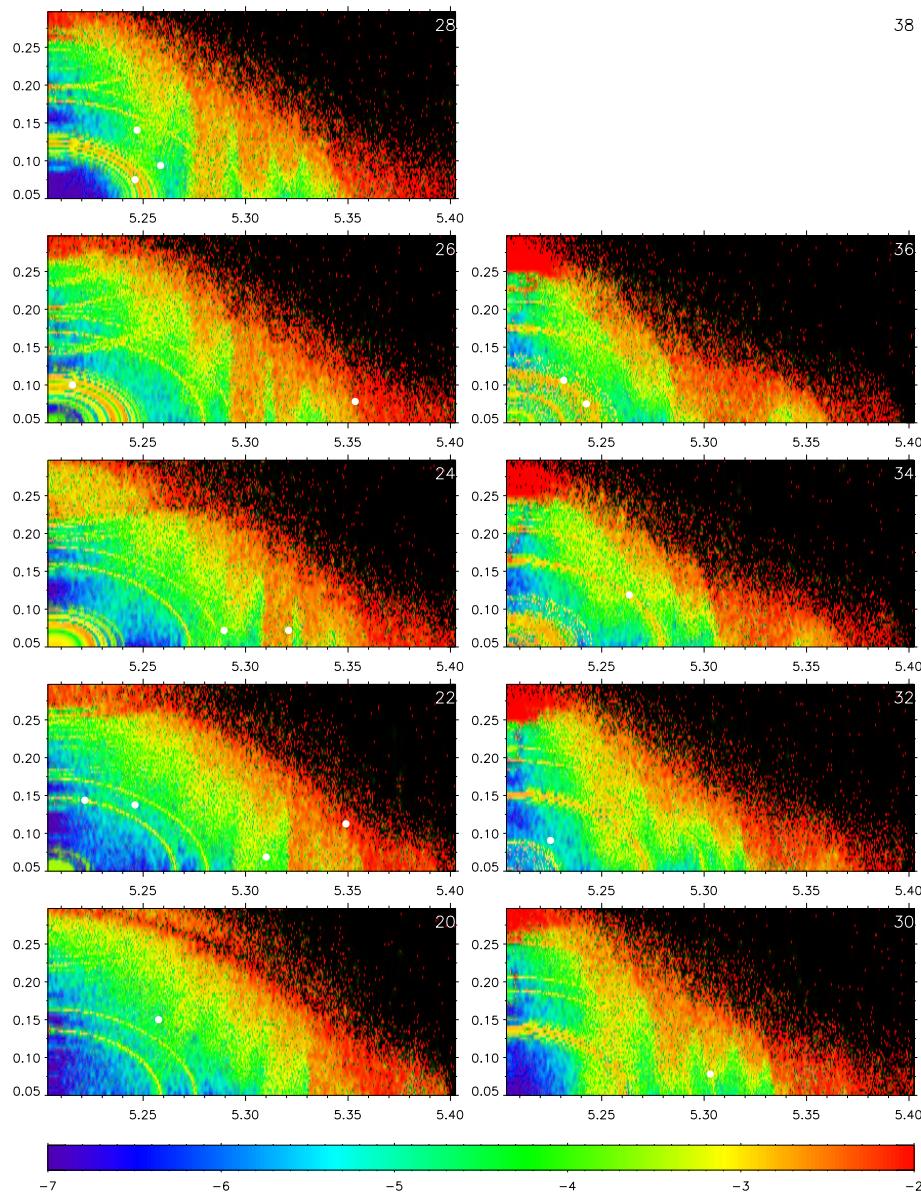
is minimal in the grid  $\tilde{\mathcal{D}}$ .

O  
B  
S  
E  
R  
V  
E  
D  
  
T  
R  
O  
J  
A  
N  
S

(I)



O  
B  
S  
E  
R  
V  
E  
D  
  
T  
R  
O  
J  
A  
N  
S  
  
(II)

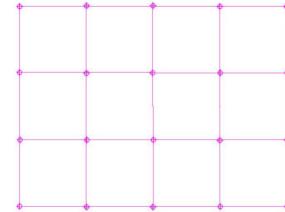


## Global dynamics: Technical details (I)

1. Initial conditions (orbital elements) of Jupiter and Saturn.

Take a grid of points in the  $(a, e)$  plane:

2.  $(a, e) \in \mathcal{A} = [5.20, 5.40] \times [0.05, 0.30]$  and  
 $I = I_5, \lambda = \lambda_5 - \frac{\pi}{3}, \varpi = \varpi - \frac{\pi}{3}, \Omega = \Omega_5$ .



3. For every initial conditions  $(a, e)$ , perform a 5 Myr. integration with the *symplectic integrator* of the family  $SABA_n$ : J. Laskar and P. Robutel, “High order symplectic integrators for perturbed Hamiltonian systems”, *Celestial Mech.*, 80:39–62, 2001.
  - (a) Take a sample of  $N$  equidistant points and do frequency analysis of this sample of the trajectory.
  - (b) Keep the first 10 frequencies (with larger amplitude).
  - (c) Choose a base for  $(\nu, g, s)$ . In our case,

$$\nu \in [7000, 9500] \quad g \in [250, 450] \quad s \in [-50, 10] \quad (\text{``/yr})$$

Periods:  $\nu \rightsquigarrow 150 \text{ yr}$        $g \rightsquigarrow 4000 \text{ yr}$        $s \rightsquigarrow 65000 \text{ yr}$

- (d) Identify  $(\nu, g, s)$  of the particular trajectory.

## Global dynamics: Technical details (II)

4. Repeat point (3.) for the interval 5 Myr. to 10 Myr.

Keep the frequencies  $(\nu', g', s')$ .

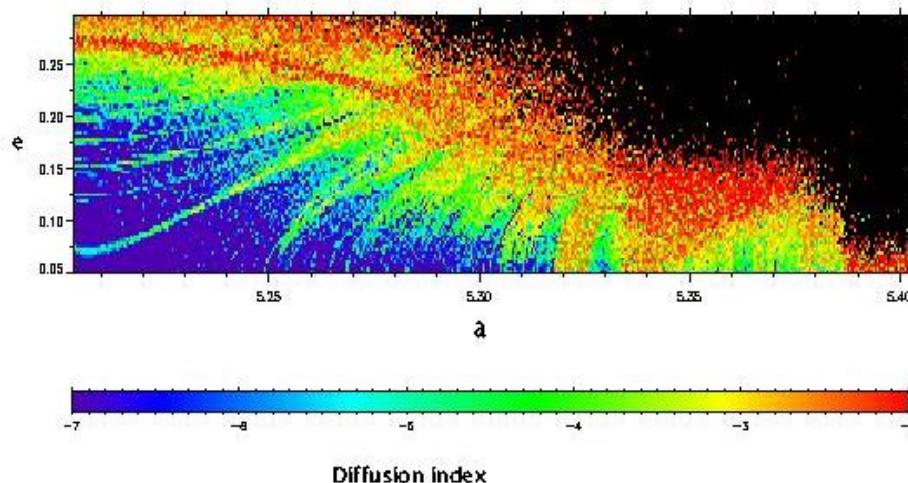
5. Compute the diffusion index  $\delta\nu = 1 - \frac{\nu'}{\nu}$

6. Save a file where every entry (line) is

$a$        $e$        $\nu$        $g$        $s$        $\delta\nu$        $\delta g$        $\delta s$

7. Perform a contour plot assigning a color to  $\delta\nu$  for every point  $(a, e)$ :

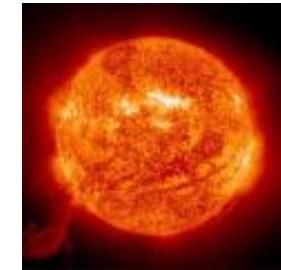
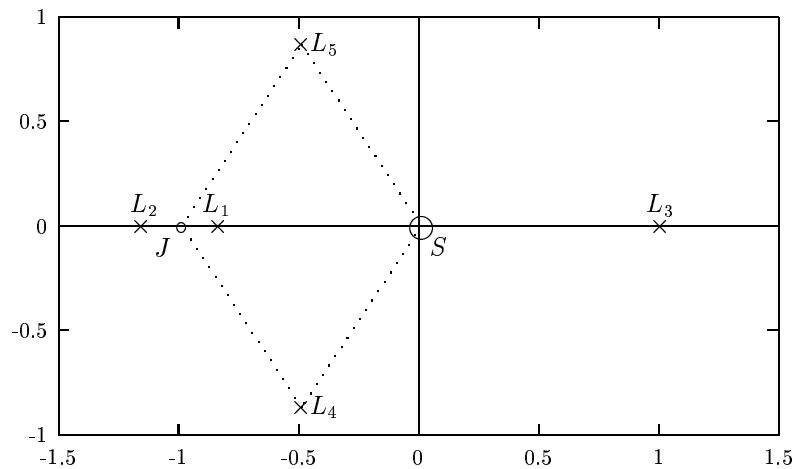
- PGPlot
- Gnuplot
- Matlab



## COMPUTER LAB



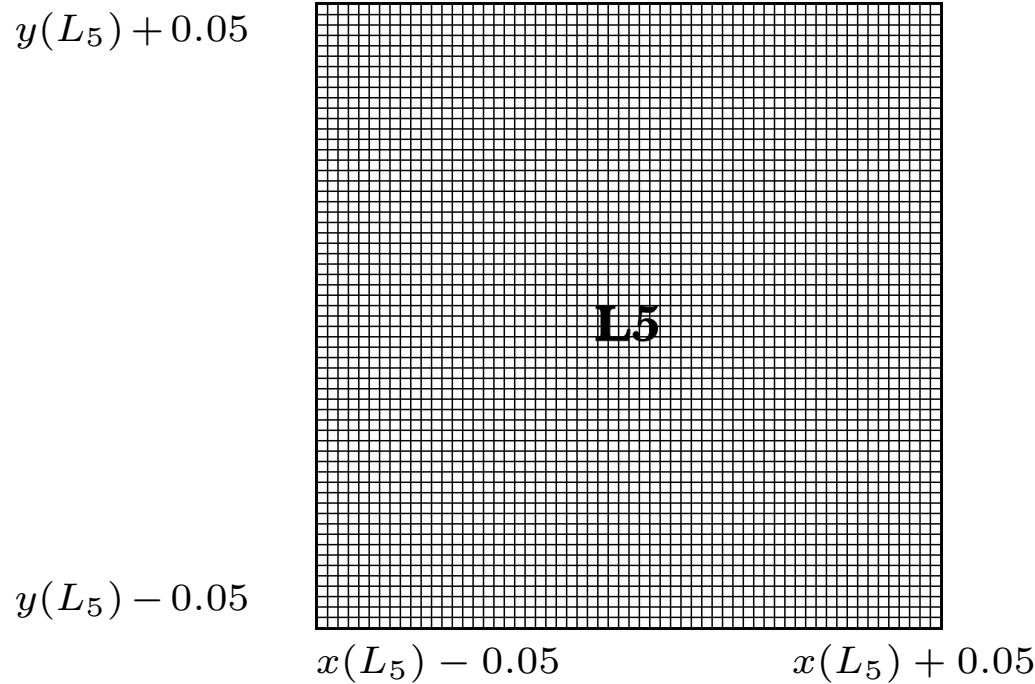
## Computer Lab: 2 dof model, planar RTBP



$$\begin{aligned} H = & \frac{1}{2} (p_x^2 + p_y^2) + (yp_x - xp_y) \\ & - \frac{1 - \mu}{\sqrt{(x - \mu)^2 + y^2}} - \frac{\mu}{\sqrt{(x - \mu + 1)^2 + y^2}} \end{aligned}$$

## Computer Lab: Initial conditions

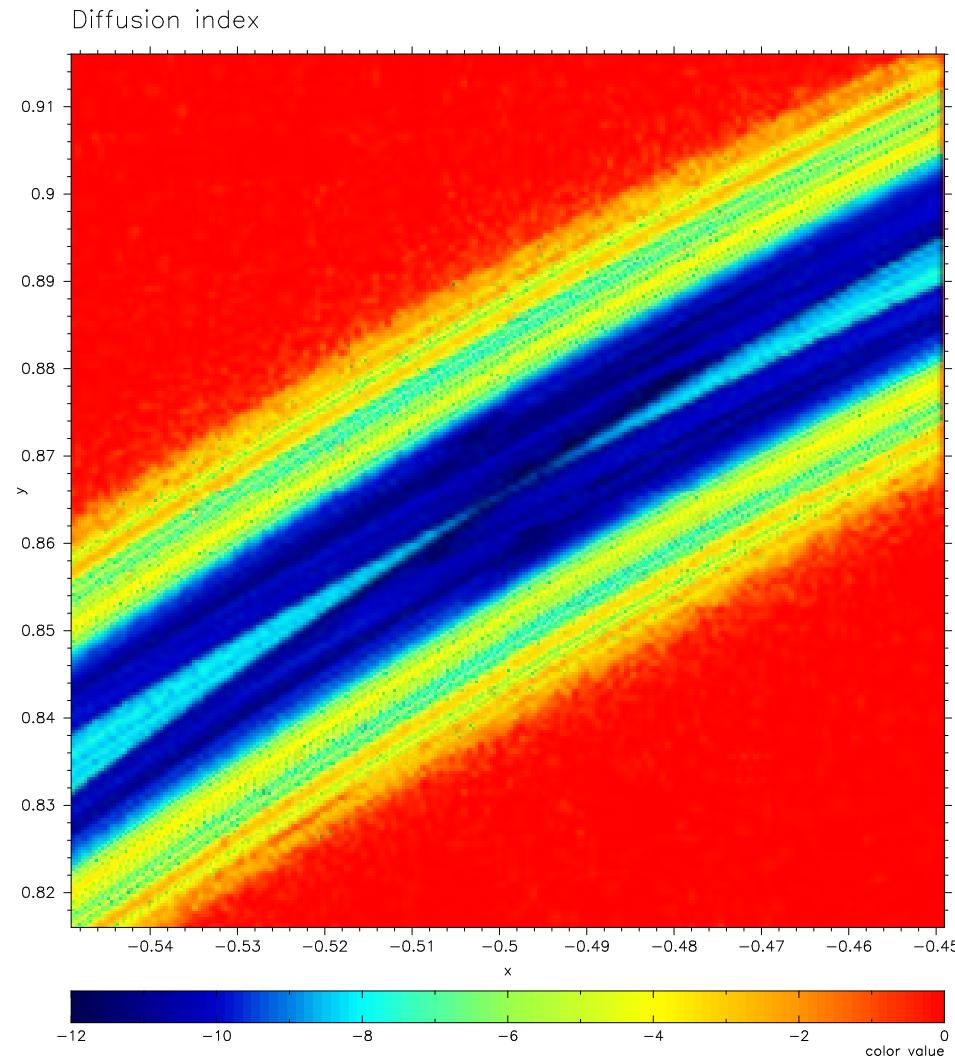
$$H(x, y, p_x, p_y) : \quad (x, y, p_x, p_y) \in \mathbb{R}^4$$



$$L_5 : \quad x(L_5) = \mu - 0.5 \quad y(L_5) = \frac{\sqrt{3}}{2} \quad p_x(L_5) = -y(L_5) \quad p_y(L_5) = x(L_5)$$

Fix:  $p_x = p_x(L_5)$  and  $p_y = p_y(L_5)$ .

## Computer Lab: Final Goal



## Computer Lab: Tools

- Numerical Integration of ODEs:
  - Taylor method.
  - Your favorite integrator.
- Furiane: Refined Fourier Analysis
- PGPLOT.

See you tomorrow at the LAB !!

