



Synchronization Patterns in Firing Rate Models with Synaptic Delay

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Introduction

- Brain oscillations, Hans Berger 1929 (EEG)
- Display a broad range of frequencies
- Correlated with sleep stages & tasks
- The reflect some coordination of spike discharges in large ensembles of neurons
- Inhibition largely involved, particularly in "fast oscillations" (>30Hz)
- Mathematical models:
 - Inhibition + Synaptic Delays

Modeling populations of neurons

• Microscopic modeling: Networks of spiking neurons

Describe the activity of a population of N neurons by a set of O(N) ordinary differential equations, coupled via a given connectivity matrix

- Macroscopic modeling: Firing rate models Wilson-Cowan, 1972
 - Activity of a neuronal **population**, single "averaged" firing rate *r(t)*
 - Heuristic (not derived from spiking neuron networks)
 - Fail to describe **Synchronization**
 - Analytically and numerically convenient, also from the point of view of experimental neuroscience

Orientation selectivity in V1





Hubel and Wiesel 1981



- Firing rate codes for features of visual stimuli
- Selectivity also: Frequency: A1, Direction of the arm: M1; Spatial location: PFC...
- Redundancy: Nearby neurons have similar response properties: Population Firing Rate

Nearby neurons display similar response properties



Blasdel, J. Neuroscience (1992)

Fast oscillations in Heuristic Firing Rate Models

$$\tau \dot{r} = -r + \Phi(-Jr(t-D) + I)$$

- r(t): Firing rate (at time t)
- $\Phi(I)$: Transfer function (f-I curve)
- -Jr(t-D): Time delayed, inhibitory synaptic current
- *I*: External currents

Linear Stability analysis

- Fixed point $r^* = \Phi(Jr^* + I)$
- Characteristic equation $\tau \lambda = -1 + J e^{-\lambda D}$
- Hopf (supercritical): $\tan(\Omega_c D) = -(\tau/D)\Omega_c D$
- $T_c = \frac{2\pi}{\Omega_c} \in (2D, 4D)$
- $D \sim 5 \text{ms} \rightarrow T_c \in (10, 20) \text{ms}$: Fast Oscillations





Roxin, Brunel, Hansel, PRL (2005); Brunel, Hakim, Chaos (2008)

Fast oscillations in spiking neuron models

• Often, neurons do not fire at the freq. of the mean field

Dichotomy between Macroscopic & Microscopic dynamics

- In contrast w. Collective Synchronization Winfree J. Theor Biol. 1967, Kuramoto 1975, 1984.
- Interesting macroscopic, oscillatory state, in spiking neuron networks:

Sparse Synchronization Brunel & Hakim 1999

Microscopic analog of fast oscillations Sparse Synchronization

- Networks of **non-oscillatory**, spiking neurons
- Strongly driven by **noise**
- Inhibition
- Synaptic delays (fixed and/or synaptic kinetics)



Micro. description: Sparse Sync Macro. description: Fast Osc. in FRM

FRM are heuristic: not *exact* derivations from spiking neuron networks

Can we find an **exact** correspondence between:

Fast Oscillations at the macroscopic level

Some different state at the microscopic level

Derivation of Exact Firing Rate Equations from a network of Spiking Neurons (with quenched noise, only!)

Montbrió, Pazó and Roxin, Physical Review X 2015

Spiking neurons Quadratic Integrate & Fire model (QIF)

The QIF model is the normal form of a SNIC bifurcation



Oscillatory dynamics





Ermentrout, Kopell, SIAM 1986

Dynamics of the QIF model

$$V = I + V^2$$
, if $V \ge V_{\text{peak}}$, then $V \leftarrow V_{\text{reset}}$

Excitable dynamics:



E. Izhikevich, "Dynamical Systems in Neuroscience", 2007

Ensemble of recurrently coupled QIF neurons with synaptic time delay

$$\tau \dot{V}_j = V_j^2 + I_j,$$
$$I_j = \eta_j + J s_D,$$

- Coupling: *J*>0: Excitation; *J*<0: Inhibition
- Mean synaptic activity ($s_D = s(t-D)$): $s_D = \frac{\tau}{N\tau_s} \sum_{i=1}^{N} \sum_{k} \int_{t-D-\tau_s}^{t-D} \delta(t'-t_j^k) dt'.$
- Fast synapses (τ_s ->0): $s_D = \tau r_D$

Time delayed, Population-Averaged Firing Rate

Thermodynamic limit Continuous formulation

$ ho(V \eta,t)dV$	Fraction of neurons with V between V and $V+dV$
	and parameter η at time <i>t</i>

 $g(\eta)$ PDF of the currents η

The Continuity Equation is

$$\partial_t \rho + \partial_V \left[(V^2 + \eta + Js + I)\rho \right] = 0$$

For each value of η !! Then the total density at time *t* is given by: $\int_{-\infty}^{\infty} \rho(V|\eta, t) g(\eta) d\eta$

Stationary solutions

 $-\partial_t \rho = \partial_V \left(\rho [V^2 + \eta] \right)$

• If
$$\eta > 0$$
: $\rho(V|\eta) = \frac{C(\eta)}{V^2 + \eta}$

• If $\eta \leq 0$: $\rho(V|\eta) = \delta(V - \tilde{C}(\eta))$

Lorentzian Ansatz

$$\rho(V|\eta) = \frac{1}{\pi} \frac{x(\eta)}{(V-y(\eta))^2 + x(\eta)^2}$$

General solutions?

• Lorentzian Ansatz:
$$\rho = \frac{1}{\pi} \frac{1}{(V-y)^2 + x^2}$$

• Continuity Eq: $-\partial_t \rho = \partial_V \left(\rho [V^2 + \eta] \right)$

We substitute the LA into the continuity eq

•
$$\partial_t \rho = \frac{1}{\pi} \frac{1}{((V-y)^2 + x^2)^2} \left(\dot{x} [(V-y)^2 + x^2] - x [2x\dot{x} - 2\dot{y}(V-y)] \right)$$

•
$$\partial_V \left(\rho [V^2 + \eta] \right) = \frac{-2(V-y)x}{\pi ((V-y)^2 + x^2)^2} [V^2 + \eta] + \frac{2Vx}{\pi ((V-y)^2 + x^2)}$$

Equating the expressions

$$-\dot{x}\left((V-y)^2 + x^2\right) + 2x\left(\dot{x} - \dot{y}(V-y)\right) = -2(V-y)x[V^2+\eta] + 2Vx\left[(V-y)^2 + x^2\right]$$

The identity must hold at all orders!!

•
$$O(V^2)$$
: $\dot{x} = 2xy$

•
$$O(V): \dot{y} = y^2 - x^2 + \eta$$

• O(1): Linear combination of previous equations

Dynamics in the Lorentzian manifold

$$\rho(V|\eta,t) = \frac{1}{\pi} \frac{x(\eta,t)}{[V-y(\eta,t)]^2 + x(\eta,t)^2} \qquad \partial_t \rho + \partial_V \left[(V^2 + \eta + Js + I)\rho \right] = 0$$
Lorentzian ansatz
Continuity equation
$$w(\eta,t) \equiv x(\eta,t) + iy(\eta,t)$$

$$\partial_t w(\eta,t) = i \left[\eta + Js(t) - w(\eta,t)^2 + I(t) \right]$$

$$s(t) = r(t) : \text{Fast Synapses}$$

Closing this equation requires to express w as a function of r and some other meaningful macroscopic observables

Lorentzian Ansatz Firing Rate & Mean Membrane potential

Firing Rate = Prob flux at threshold: $r(\eta, t) = \rho(V \to \infty | \eta, t) \dot{V}(V \to \infty | \eta, t)$

Firing Rate

$$x(\eta, t) = \pi r(\eta, t) \qquad r(t) = \frac{1}{\pi} \int_{-\infty}^{\infty} x(\eta, t) g(\eta) d\eta$$

Mean Membrane potential

$$y(\eta, t) = \text{P.V.} \int_{-\infty}^{\infty} \rho(V|\eta, t) V \, dV$$

$$v(t) = \int_{-\infty}^{\infty} y(\eta, t) g(\eta) d\eta$$

Firing Rate Model for QIF neurons

Lorentzian distribution of currents

$$g(\eta) = \frac{1}{\pi} \frac{\Delta}{(\eta - \bar{\eta})^2 + \Delta^2}$$

Cauchy Residue's theorem to solve Ott & Antonsen, *Chaos*, 2008

$$\begin{aligned} r(t) &= \frac{1}{\pi} \int_{-\infty}^{\infty} x(\eta, t) g(\eta) d\eta \\ v(t) &= \int_{-\infty}^{\infty} y(\eta, t) g(\eta) d\eta \end{aligned}$$

$$egin{aligned} & au \dot{r} = rac{\Delta}{\pi au} + 2rv, \ & au \dot{v} = v^2 + ar{\eta} + J au r_D - au^2 \pi^2 r^2, \end{aligned}$$

Analysis of the FRE without delay



Phase portrait: Movie!



Comparison Spiking Neurons vs. Firing Rate Eqs.



Firing Rate Equations with time delays

Pazó & Montbrió, Physical Review Letters 2016

Linear Stability Analysis of Incoherence

$$egin{aligned} & au \dot{r} = rac{\Delta}{\pi au} + 2rv, \ & au \dot{v} = v^2 + ar{\eta} + J au r_D - au^2 \pi^2 r^2, \end{aligned}$$

 $\tau=\eta=1$, without loss of generality

For identical neurons, the only fixed point is:

$$\left((J+\sqrt{J^2+4\pi^2})/(2\pi^2),0\right)$$

Incoherent state (splay state)

Linearizing around the f.p. and imposing the cond. of marginal stab: λ = i Ω

Hopf boundaries:

$$\Omega_n = n\pi/D.$$

$$J_{H}^{(n)} = \pi (\Omega_{n}^{2} - 4) \times \begin{cases} (6\Omega_{n}^{2} + 12)^{-1/2} & \text{for odd } n \\ (2\Omega_{n}^{2} - 4)^{-1/2} & \text{for even } n \end{cases}$$

Incoherence (Hopf) boundaries



Weakly nonlinear analysis (two timing)



Synchronization boundaries

(obtained analytically via linear stability analysis of sync)



Phase diagram



Fast oscillations in FRM T=2D1 D 0 15 0 5 Delayed Inhibition (D) After a new time D prevents an increase of activity there is no inhibition... Firing is possible again!

Micro vs. Macro dynamics



QUASIPERIODIC PARTIAL SYNCHRONIZATION in inhibitory networks

Van Vresswijk, Phys Rev E (1996)

Mohanti, Politi, J. Phys A (2006); Rosenblum, Pikovsky, PRL 2007; Pikovsky, Rosenblum, Physica D (2009); Olmi, Politi, Torcini, EPL (2010); Luccioli et al, Phys Rev Lett (2012); Politi, Rosenblum, PRE (2015)...

period of QPS is always 2D



Symmetry of limit cycle implies T=2D



Period of oscillations remains constant

Limit cycle is symmetric $V \rightarrow -V$

Using FRE for Δ =0, this symmetry implies that:

$$T_m = \frac{2D}{m}$$
, with $m = 1, 3, \dots$

The symmetry is broken at period doubling bif...

Transition from QPS to Macroscopic Chaos



Neurons are not chaotic, though...

Collective chaos

Transition from QPS to Collective Chaos period-doubling cascade

Using the FREs we find:



In the thermodynamic limit the system shows genuine Collective Chaos

Onset of QPS and (weak) heterogeneity





TC bifs. are not robust

bistability remains though!

Macroscopic chaos in heterogeneous networks



Summary

- Exact derivation of FRE for networks of QIF neurons
- Using the FRE we related
 - Quasiperiodic Partial Synchronization
 - Collective Chaos

with (non-trivial) Fast Oscillations in Inhibitory Networks

- Transition from QPS to CC via period doubling cascade
- Fast Oscillations in Inhibitory networks may correspond to Many possible Micro-states

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