

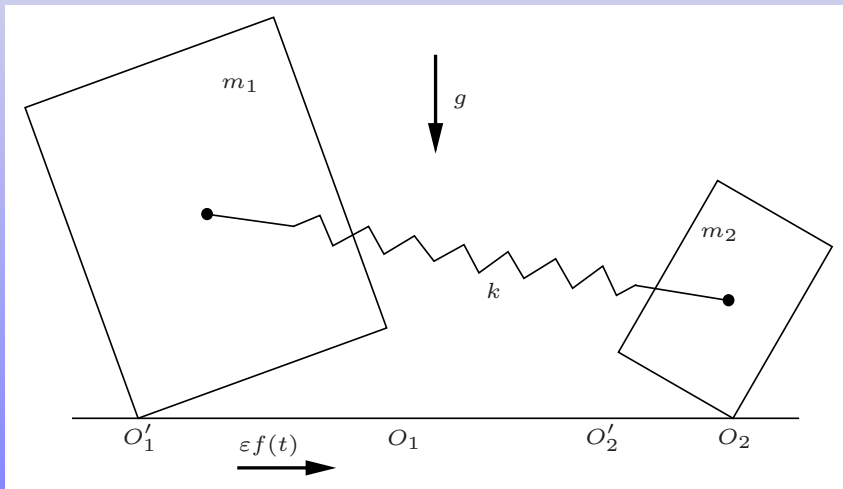
On energy accumulation in a class of piecewise-defined Hamiltonian systems

Albert Granados, John Hogan, Tere Seara

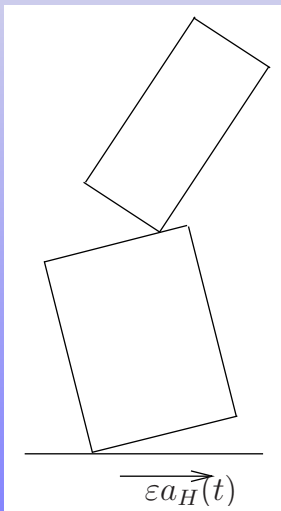
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Motivation



Motivation

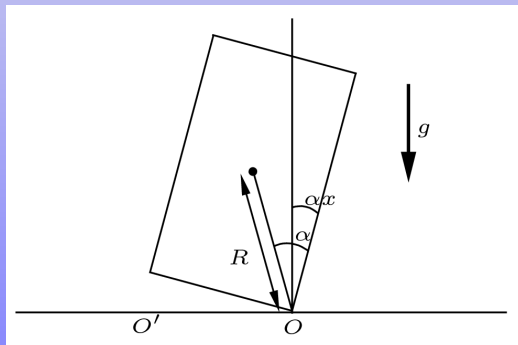


- 1 Dynamics of the rocking block
 - Phase portrait
 - Generalization: piecewise-defined “Hamiltonian”
- 2 The coupled system
 - System description
 - The unperturbed case
 - The perturbed case
- 3 The scattering map
 - Intersection of the heteroclinic manifolds
 - Properties of the scattering map
- 4 Numerical results

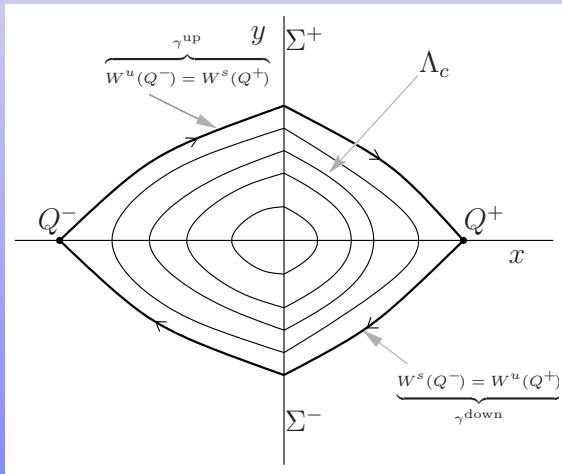
The rocking block

$$\alpha \ddot{x} + \sin(\alpha(1-x)) = 0 \quad \text{if } x > 0$$

$$\alpha \ddot{x} - \sin(\alpha(1+x)) = 0 \quad \text{if } x < 0$$



Phase portrait



Piecewise-defined system

Let us split the plane in two sets

$$S^+ = \{(x, y) \in \mathbb{R}^2 \mid x > 0\}$$

$$S^- = \{(x, y) \in \mathbb{R}^2 \mid x < 0\}$$

separated by the switching manifold

$$\Sigma = \Sigma^+ \cup \Sigma^- \cup (0, 0)$$

where

$$\Sigma^+ = \{(x, y) \in \mathbb{R}^2 \mid x = 0, y > 0\}$$

$$\Sigma^- = \{(x, y) \in \mathbb{R}^2 \mid x = 0, y < 0\}.$$

Piecewise-defined system

We consider the piecewise smooth system

$$\begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} = \begin{cases} \mathcal{X}_0^+(x, y) & \text{if } (x, y) \in S^+ \cup \Sigma^+ \\ \mathcal{X}_0^-(x, y) & \text{if } (x, y) \in S^- \cup \Sigma^- \end{cases}$$

which we assume “Hamiltonian”

Piecewise-defined system

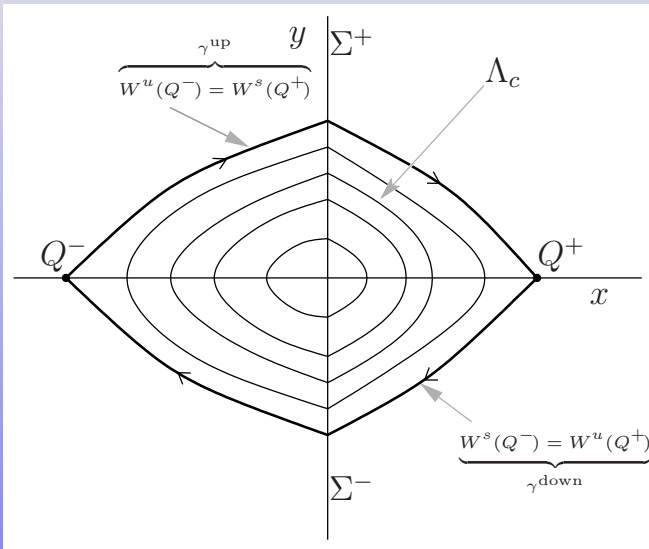
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which we assume "Hamiltonian", with Hamiltonian

$$X(x, y) := \frac{y^2}{2} + Y(x) := \begin{cases} \frac{y^2}{2} + Y^+(x) & \text{if } (x, y) \in S^+ \cup \Sigma^+ \\ \frac{y^2}{2} + Y^-(x) & \text{if } (x, y) \in S^- \cup \Sigma^- \end{cases}$$

with $Y^\pm \in C^\infty(\mathbb{R})$ satisfying $Y^+(0) = Y^-(0)$ but $(Y^+)'(0) \neq (Y^-)'(0)$.



The coupled system

We consider

$$H_\varepsilon(x, y, u, v, t) := X(x, y) + U(u, v) + \varepsilon h(x, y, u, v, t), \quad \varepsilon > 0,$$

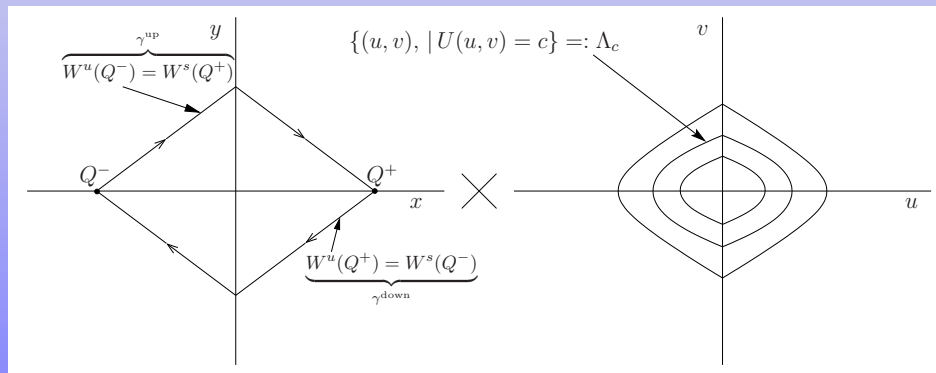
with h T -periodic in t :

$$h(x, y, u, v, t) = h(x, y, u, v, t + T)$$

and X and U are piecewise-defined Hamiltonian modelling two rocking blocks.

Invariant objects for the unperturbed case

If $\varepsilon = 0$, the coupled system is a cross product of two rocking blocks $\times \mathbb{T}$



Invariant objects for the unperturbed case

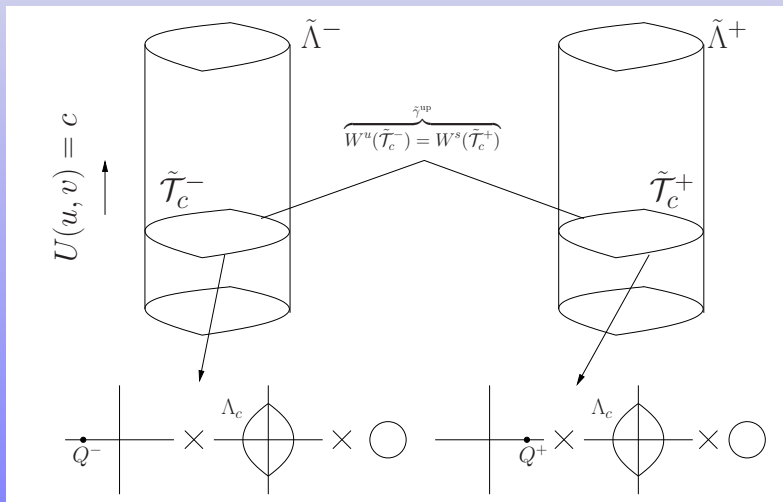
- $\tilde{\mathcal{T}}_c^\pm = Q^+ \times \Lambda_c \times \mathbb{T} =$
 $\{(x, y, u, v, s) \mid (x, y) = Q^\pm, U(u, v) = c, s \in \mathbb{T}\}$

Invariant objects for the unperturbed case

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 $\{(x, y, u, v, s) \mid (x, y) = Q^\pm, U(u, v) = c, s \in \mathbb{T}\}$
- $W^s(\tilde{\mathcal{T}}_c^+) = W^u(\tilde{\mathcal{T}}_c^-)$
 $= W^s(Q^+) \times \Lambda_c \times \mathbb{T} = W^u(Q^-) \times \Lambda_c \times \mathbb{T}$

Invariant objects for the unperturbed case

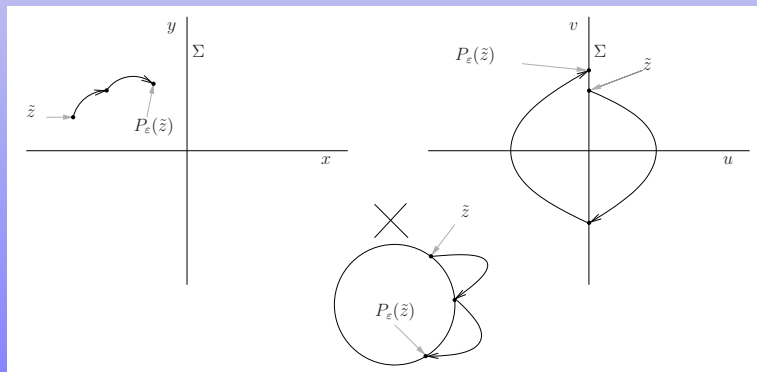
- $\tilde{\mathcal{T}}_c^\pm = Q^\pm \times \Lambda_c \times \mathbb{T} = \{(x, y, u, v, s) \mid (x, y) = Q^\pm, U(u, v) = c, s \in \mathbb{T}\}$
- $W^s(\tilde{\mathcal{T}}_c^+) = W^u(\tilde{\mathcal{T}}_c^-)$
 $= W^s(Q^+) \times \Lambda_c \times \mathbb{T} = W^u(Q^-) \times \Lambda_c \times \mathbb{T}$
- $\tilde{\Lambda}^- = \bigcup_{c \in [c_1, c_2]} \tilde{\mathcal{T}}_c^- = \bigcup_{c \in [c_1, c_2]} Q^- \times \Lambda_c \times \mathbb{T}$
- $\tilde{\Lambda}^+ = \bigcup_{c \in [c_1, c_2]} \tilde{\mathcal{T}}_c^+ = \bigcup_{c \in [c_1, c_2]} Q^+ \times \Lambda_c \times \mathbb{T}$

C^0 Invariant manifolds

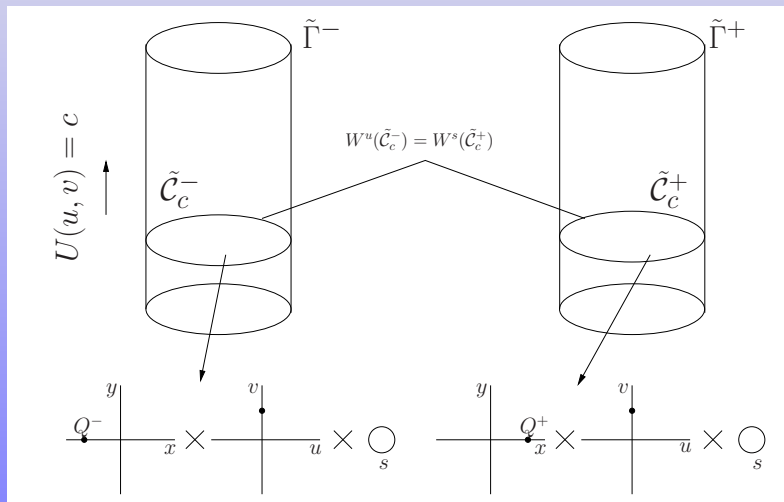
Persistence when $\varepsilon > 0$

We don't have theory that provides the persistence of these invariant manifolds because of the non-smoothness. Instead we use the impact map:

$$P_\varepsilon : \mathbb{R}^2 \times \Sigma \times \mathbb{T} \longrightarrow \mathbb{R}^2 \times \Sigma \times \mathbb{T}$$



NHIM's for the impact map



Perturbed invariant manifolds for the map

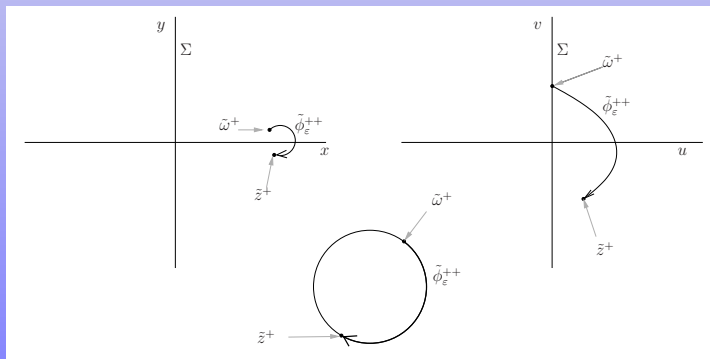
Using that P_ε is smooth, we can apply classical theory of normally hyperbolic manifolds to have the existence of

$$\tilde{\Gamma}_\varepsilon^\pm, W^{u,s}(\tilde{\Gamma}_\varepsilon^\pm)$$

Perturbed invariant manifolds for the flow

$$\tilde{\omega}^+ = (\omega^+, s^+) \in \tilde{\Gamma}_\varepsilon^+ \longrightarrow \tilde{z} \in \tilde{\Lambda}_\varepsilon^+$$

$$\tilde{z}^+ = (z^+, s^+ + \tau) = \tilde{\phi}(\tau; \tilde{\omega}^+; \varepsilon), \quad 0 \leq \tau < \Pi_s (P_\varepsilon(\tilde{\omega}^+) - s^+)$$



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$$\tilde{\omega}^+ = (\omega^+, s^+) \in \tilde{\Gamma}_\varepsilon^+$$

$$\tilde{\omega}^s = (\omega^s, s^s) \in W^s(\tilde{\Gamma}_\varepsilon^+)$$

$$|P_\varepsilon^n(\tilde{\omega}^+) - P_\varepsilon^n(\tilde{\omega}^s)| \longrightarrow 0$$

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$$\tilde{z}^+ = \tilde{\phi}(\tau; \tilde{\omega}^+; \varepsilon) \in \tilde{\Lambda}_\varepsilon^+$$

$$\tilde{z}^s = \tilde{\phi}(\tau; \tilde{\omega}^s; \varepsilon) \notin W^s(\tilde{z}^+)$$

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$$\tau' = s^+ - s^s$$

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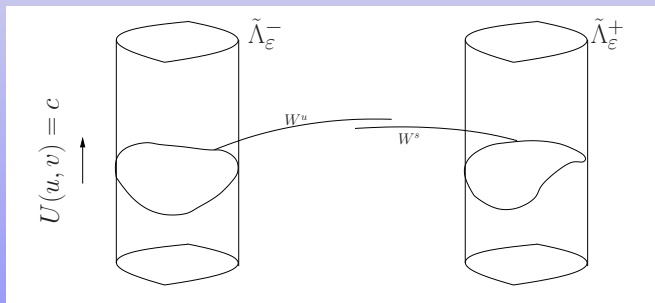
$$\tilde{z}^s = \tilde{\phi}(\tau + \tau'; \tilde{\omega}^s; \varepsilon) \in W^s(\tilde{z}^+)$$

$$\tau' = s^+ - s^s$$

$$|\tilde{\phi}(t; \tilde{z}^+; \varepsilon) - \tilde{\phi}(t; \tilde{z}^s; \varepsilon)| \longrightarrow 0$$

Intersection of the heteroclinic manifolds

When $\varepsilon > 0$ there exist 3-dimensional invariant manifolds

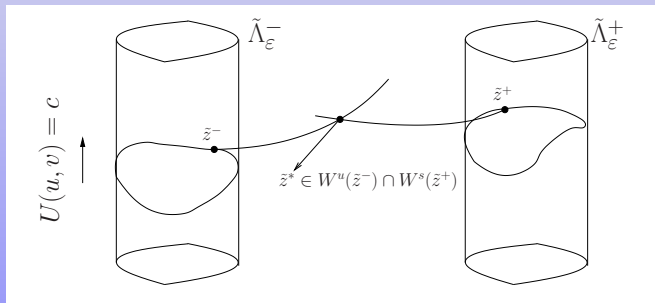


with 4-dimensional $W^u(\tilde{\Lambda}_\varepsilon^-)$ and $W^s(\tilde{\Lambda}_\varepsilon^+)$

Do they intersect??

Intersection of the heteroclinic manifolds

Yes, under certain conditions



$W^u(\tilde{\Lambda}_\epsilon^-) \cap W^s(\tilde{\Lambda}_\epsilon^+)$ is 3-dimensional

Intersection of the heteroclinic manifolds

A sufficient condition to have (transversal) intersection of the heteroclinic manifolds is that the Melnikov function

$$\zeta \mapsto M(\zeta, \tau, v, s),$$

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In the smooth case,

$$M(\zeta, \tau, v, s) := \int_{-\infty}^{\infty} \{X, h\} \left(\phi \left(t; \underbrace{(\sigma(\zeta), \phi_U(\tau; 0, v), s)}_{\tilde{z}_0}; 0 \right) \right) dt$$

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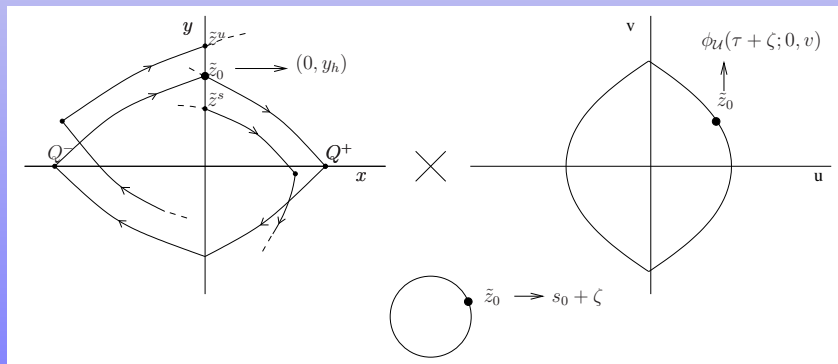
In our case,

$$M(\zeta, \tau, v, s) := \int_{-\infty}^{\infty} \{X, h\} \left(\phi \left(t; \underbrace{(0, y_h, \phi_{\mathcal{U}}(\tau + \zeta; 0, v), s + \zeta)}_{\tilde{z}_0}; 0 \right) \right) dt$$

Intersection of the heteroclinic manifolds

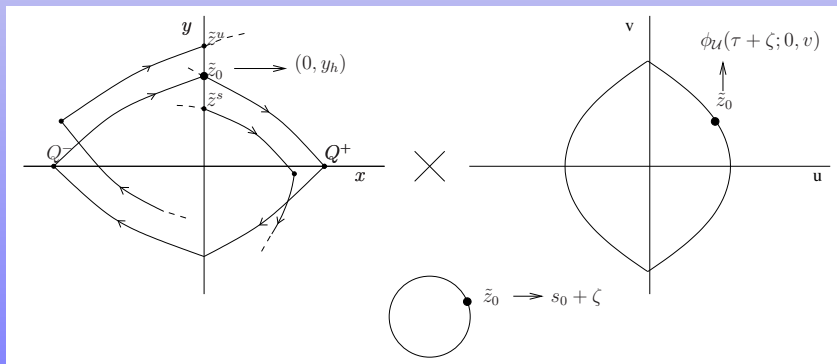
$$\tilde{z}_0 = (0, y_h, \phi_U(\tau + \zeta; 0, v), s + \zeta)$$

$$\tilde{N} = \{\tilde{z}_0 + l(0, 1, 0, 0, 0), l \in \mathbb{R}\}$$



Intersection of the heteroclinic manifolds

$$\Delta(\zeta, \tau, v, s; \varepsilon) = X(\tilde{z}^u) - X(\tilde{z}^s) = \varepsilon M(\zeta, \tau, v, s) + O(\varepsilon^2) = 0$$



Intersection of the heteroclinic manifolds

Let $\bar{\zeta}$ be a simple zero of

$$\zeta \mapsto M(\zeta, \tau, v, s)$$

$$\implies \exists \tilde{z}_0^* \in W^u(\tilde{\Lambda}_\varepsilon^-) \cap W^s(\tilde{\Lambda}_\varepsilon^+) \cap \{x = 0\}, \tilde{z}^\pm \in \tilde{\Lambda}_\varepsilon^\pm$$

$$\tilde{z}_0^* = (0, y_h, \phi_{\mathcal{U}}(\tau + \bar{\zeta}; 0, v), s + \bar{\zeta}) + O(\varepsilon)$$

$$\tilde{z}^\pm = (Q^\pm, \phi_{\mathcal{U}}(\tau + \bar{\zeta}; 0, v), s + \bar{\zeta}) + O(\varepsilon)$$

$$\left| \tilde{\phi}(t; \tilde{z}_0^*; \varepsilon) - \tilde{\phi}(t; \tilde{z}^\pm; \varepsilon) \right| \rightarrow 0, t \rightarrow \pm\infty$$

Properties of the scattering map

The scattering map becomes

$$S_\varepsilon : \quad \begin{array}{ccc} \tilde{\Lambda}_\varepsilon^- & \longrightarrow & \tilde{\Lambda}_\varepsilon^+ \\ \tilde{z}^-(\tau, v, s) & \longmapsto & \tilde{z}^+(\tau, v, s) \end{array}$$

We want to measure the difference of the energy of the points \tilde{z}^+ and \tilde{z}^- .

We use the unperturbed Hamiltonians

$$\begin{aligned} X(\tilde{z}^+) + U(\tilde{z}^+) - (X(\tilde{z}^-) + U(\tilde{z}^-)) \\ = \underbrace{X(\tilde{z}^+) - X(\tilde{z}^-)}_{O(\varepsilon^2)} + \underbrace{U(\tilde{z}^+) - U(\tilde{z}^-)}_{\Delta U_\varepsilon + O(\varepsilon^{1+\rho})} \end{aligned}$$

Properties of the scattering map

$$U(\tilde{z}^+) - U(\tilde{z}^-) = \varepsilon \Delta U + O(\varepsilon^{1+\rho})$$

$$\begin{aligned} \Delta U(\tau, v, s) = & \int_{-\infty}^0 \left(\{U, h\} (\text{unperturbed heteroclinic}) \right. \\ & \left. - \{U, h\} \left(\text{asymptotic solution in } \tilde{\Lambda}^- \right) \right) dt \\ & + \int_0^{+\infty} \left(\{U, h\} (\text{unperturbed heteroclinic}) \right. \\ & \left. - \{U, h\} \left(\text{asymptotic solution in } \tilde{\Lambda}^+ \right) \right) dt \end{aligned}$$

Properties of the scattering map

$$\lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \left(U \left(\tilde{\phi} (t; \tilde{z}^+; \varepsilon) \right) - U \left(\tilde{\phi} (-t; \tilde{z}^-; \varepsilon) \right) \right) dt$$

Properties of the scattering map

$$\begin{aligned} & \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \left(U \left(\tilde{\phi}(t; \tilde{z}^+; \varepsilon) \right) - U \left(\tilde{\phi}(-t; \tilde{z}^-; \varepsilon) \right) \right) dt \\ &= \varepsilon \left(\lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \left(\int_{-t}^t \{U, h\}(\sigma(-\bar{\zeta} + r), \phi_U(\tau + r; 0, v), s + r) dr \right) dt \right) \\ &+ O(\varepsilon^{1+\rho}) \end{aligned}$$

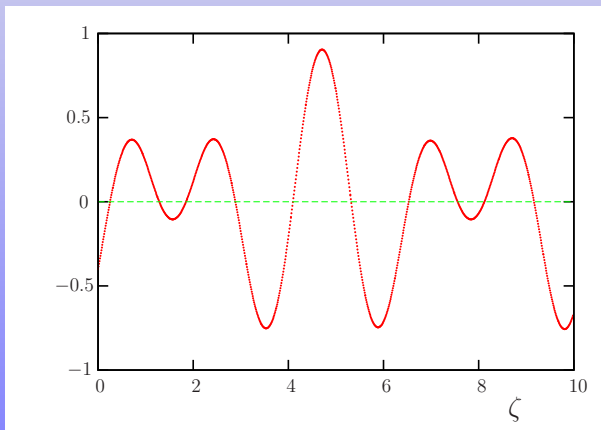
System equations

$$\begin{cases} \dot{x} = y \\ \dot{y} = x - \text{sign}(x) \\ \quad + \varepsilon(k(u - x) - \delta \cos(s)) \\ \dot{u} = v \\ \dot{v} = u - \text{sign}(u) \\ \quad + \varepsilon(k(x - u) - \delta \cos(s)) \\ \dot{s} = 1, \end{cases}$$

$$H(x, y, u, v, s) = \frac{y^2}{2} - \frac{x}{2} + |x| + \frac{v^2}{2} - \frac{u}{2} + |u| \\ + \varepsilon \left(k \left(\frac{u^2}{2} + \frac{x^2}{2} - ux \right) + \delta(x + u) \cos(s) \right),$$

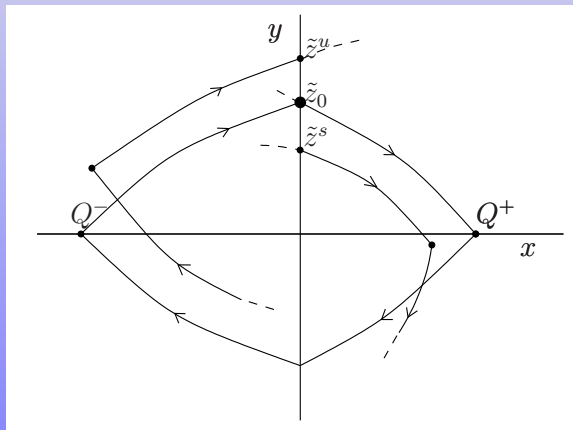
Melnikov function

Let us fix $v = 0.5$ and $\tau = s = 0$.



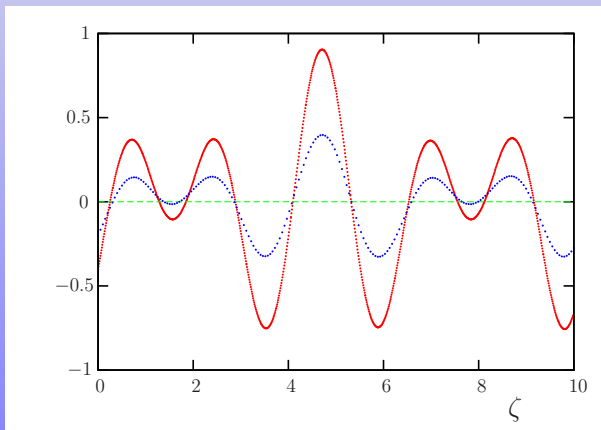
Real distance

Recalling that $\Delta(\zeta, \tau, v, s; \varepsilon) = X(\tilde{z}^u) - X(\tilde{z}^s)$.



Real distance

we get, for $\varepsilon = 0.01$ ($v = 0.5$ and $\tau = s = 0$)



$$\Delta \langle \rangle = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \left(\int_{-t}^t \{U, h\} (\sigma(-\bar{\zeta} + r), \phi_U(\tau + r; 0, v), s + r) dr \right) dt$$

Third zero $\rightarrow \Delta \langle \rangle \simeq 0.4$

Fourth zero $\rightarrow \Delta \langle \rangle \simeq -0.3$

