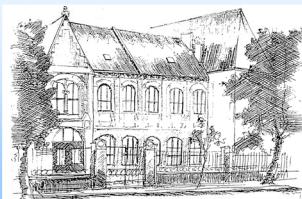


Delay equations with engineering applications

Gábor Stépán

Department of Applied Mechanics
Budapest University of Technology and Economics



Contents

Delay equations arise in mechanical systems...

... by the information system (of control), and
by the contact of bodies.

- **Linear stability & subcritical Hopf bifurcations**

- Force control and balancing – human and robotic

- Contact problems

Shimmying wheels (of trucks and motorcycles)

Machine tool vibrations

Main references

- Inspurger T, Stepan G, Stability chart for the delayed Mathieu equation, *Proceedings of the Royal Society London A* **458** (2002) 1989-1998.
- Inspurger T, Stepan G, *Semi-discretization for time-delay systems – Engineering applications*, to appear, Springer, New York, 2011.
- Inspurger T, Stepan G, Updated semi-discretization method for periodic delay-differential equations with discrete delay, *International Journal for Numerical Methods in Engineering* **61** (2004) 117-141.
- Stepan G, Inspurger T, Szalai R, Delay, parametric excitation, and the nonlinear dynamics of cutting processes, *International Journal of Bifurcation and Chaos* **15** (2005) 2783-2798.
- Orosz G, Stepan G, Subcritical Hopf bifurcations in a car-following model with reaction-time delay, *Proceedings of the Royal Society London A* **462** (2006) 2643-2670.

Non-autonomous linear RFDEs

$$M\ddot{x}(t) + \int_{-h}^0 d_\vartheta B(t, \vartheta) \dot{x}(t + \vartheta) + \int_{-h}^0 d_\vartheta K(t, \vartheta) x(t + \vartheta) = 0$$

Time-periodic systems: $B(t+T, \vartheta) = B(t, \vartheta)$

Trial solution: $x(t) = p(t)e^{\lambda t}$ $K(t+T, \vartheta) = K(t, \vartheta)$

$$p(t+T) = p(t) = \sum_{k=0}^{+\infty} (A_k \cos(k \frac{2\pi}{T} t) + B_k \sin(k \frac{2\pi}{T} t))$$

Hill's infinite dimensional determinant \Rightarrow

characteristic function \Rightarrow characteristic roots λ

$\operatorname{Re} \lambda_j < 0, j=1,2,\dots \Leftrightarrow$ stability $\Leftrightarrow |\mu_j| < 1, j=1,2,\dots$

for characteristic multipliers $\mu = e^{\lambda T}$ of fund. op. at T

The delayed Mathieu equation

Analytically constructed stability chart for testing numerical methods and algorithms

$$\ddot{x}(t) + (\delta + \varepsilon \cos t)x(t) = b x(t - 2\pi)$$

Time delay and time periodicity are equal:

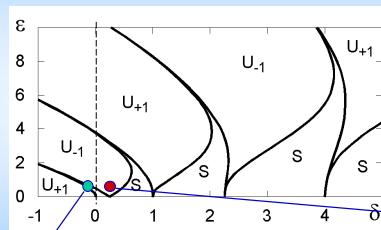
$$T = \tau = 2\pi$$

$b = 0$ Mathieu equation (1868)

$\varepsilon = 0$ Delayed oscillator (1941)

Stability chart – Mathieu equation

$$\ddot{x}(t) + (\delta + \varepsilon \cos t)x(t) = 0$$



Floquet (1883)

Hill (1886)

Rayleigh(1887)

van der Pol & Strutt (1928)

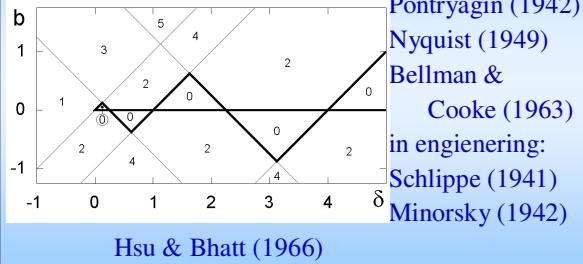
Swing (2000BC)

Strutt – Ince diagram (1956)

Stephenson (1908), Swinney (2004), Zelei (2005)

Stability chart – delayed oscillator

$$\ddot{x}(t) + \delta x(t) = b x(t - 2\pi)$$



The delayed Mathieu equation

$$\ddot{x}(t) + (\delta + \varepsilon \cos t)x(t) = b x(t - 2\pi)$$

$$x(t) = \sum_{k=0}^{\infty} (A_k e^{ikt} + B_k e^{-ikt}) e^{\lambda t} + \sum_{k=0}^{\infty} (\bar{A}_k e^{-ikt} + \bar{B}_k e^{ikt}) e^{\bar{\lambda} t}$$

$$x(t) = \sum_{k=-\infty}^{\infty} C_k e^{(\lambda+ik)t} + \bar{C}_k e^{(\bar{\lambda}-ik)t}$$

Harmonic balance
⇒ Hill's determinant

$$\det \begin{pmatrix} \ddots & \ddots & \ddots & 0 & 0 & \dots \\ \dots & 0 & \frac{\varepsilon}{2} & \delta + (\lambda + ik)^2 - b e^{-2\pi(\lambda + ik)} & \frac{\varepsilon}{2} & 0 & \dots \\ \dots & 0 & \ddots & \ddots & \ddots & \ddots & \ddots \end{pmatrix} = 0$$

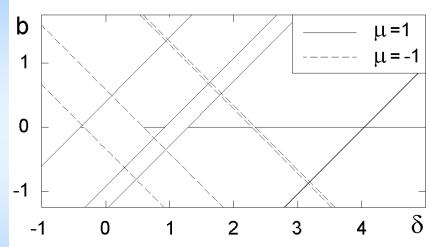
Stability boundaries

$$\lambda = \pm i\omega \Rightarrow \delta - (\omega + k)^2 - b \cos(2\pi\omega) \pm i b \sin(2\pi\omega)$$

In the main diagonal, there are 4 cases to check if 0:

- 1) $\omega \neq j/2, b \neq 0$ then the main diagonal is complex, the Hill's determinant is non-zero
- 2) $b = 0$ then there are possible stability boundaries
- 3) $b \neq 0, \omega = j/2 = h \Rightarrow \delta - b - (k + j/2)^2$ then $D(i\omega, \delta, b, \varepsilon) = f(\delta - b, \varepsilon)$ and $\mu = e^{ih\pi} = 1$ (SN)
- 4) $b \neq 0, \omega = j/2 = (2h+1)/2 \Rightarrow \delta + b - (k + j/2)^2$ then $D(i\omega, \delta, b, \varepsilon) = f(\delta + b, \varepsilon)$ and $\mu = e^{i(2h+1)\pi} = -1$ (flip)

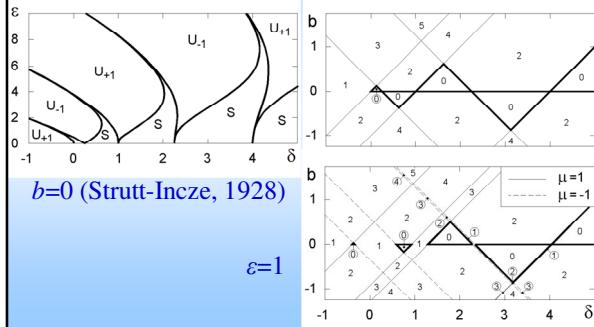
Stability boundaries remain lines in (δ, b)



$$\text{To select stable ones: } \frac{\partial \operatorname{Re} \lambda}{\partial b} = \frac{2\pi \Gamma_N^2}{(2\pi b(-1)^j \Gamma_N)^2 + \Omega_N^2} b$$

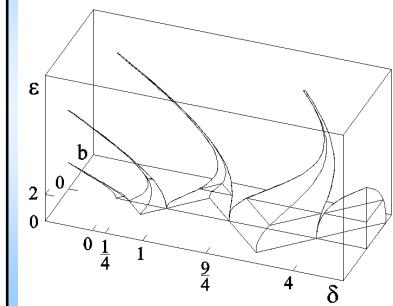
The delayed Mathieu – stability charts

$$\ddot{x}(t) + (\delta + \varepsilon \cos t)x(t) = b x(t - 2\pi)$$



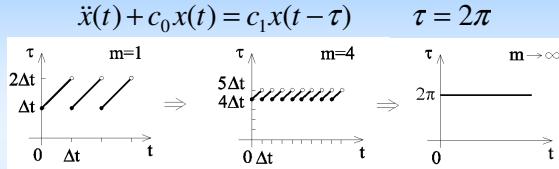
Stability chart of delayed Mathieu

$$\ddot{x}(t) + (\delta + \varepsilon \cos t)x(t) = b x(t - 2\pi)$$



Insperger,
Stépán (2002)

Semi-discretization method – introduction



The approximating DDE is *non-autonomous*

$$\ddot{x}(t) + c_0 x(t) = c_1 x(t - \tau(t)), \quad \tau(t) = t + (m - \text{int}(t / \Delta t))\Delta t$$

$$t \in [t_i, t_{i+1}] = [i\Delta t, (i+1)\Delta t] \quad \Delta t = 2\pi/(m+1/2)$$

$$\Rightarrow x(t - \tau(t)) \equiv x((i-m)\Delta t) = x_{i-m}$$

Introduction to SDM – delayed oscillator

$$\ddot{x}(t) + c_0 x(t) = c_1 x_{i-m}$$

$$x(t_i) = x_i$$

$$\dot{x}(t_i) = \dot{x}_i$$

$$x(t) = K_{1i} \cos(\sqrt{c_0}t) + K_{2i} \sin(\sqrt{c_0}t) + c_1 x_{i-m} / c_0$$

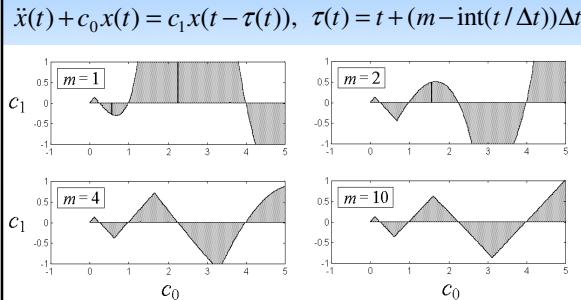
$$\dot{x}(t) = -K_{1i} \sqrt{c_0} \sin(\sqrt{c_0}t) + K_{2i} \sqrt{c_0} \cos(\sqrt{c_0}t)$$

$$x_{i+1} = a_{00} x_i + a_{01} \dot{x}_i + a_{0m} x_{i-m} \quad \mathbf{y}_i = \text{col}(\dot{x}_i, x_i, x_{i-1}, \dots, x_{i-m})$$

$$\dot{x}_{i+1} = a_{10} x_i + a_{11} \dot{x}_i + a_{1m} x_{i-m} \quad \boxed{\mathbf{y}_{i+1} = \mathbf{A} \mathbf{y}_i}$$

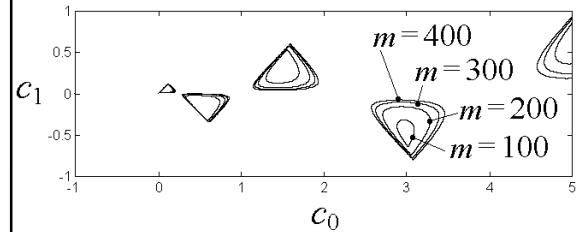
$$\det(\mu \mathbf{I} - \mathbf{A}) = 0 \Rightarrow |\mu_{1,2,\dots,m+2}| < 1 \Leftrightarrow \text{stability}$$

Delayed oscillator – stability chart by SDM

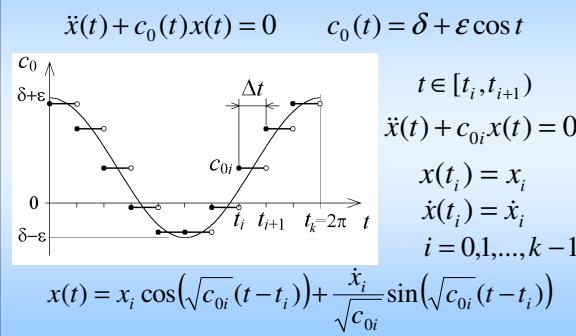


Full discretization - comparison

Discretization also w.r.t. time derivatives
– slow convergence



Introduction to SDM – Mathieu equation



SDM for Mathieu equation

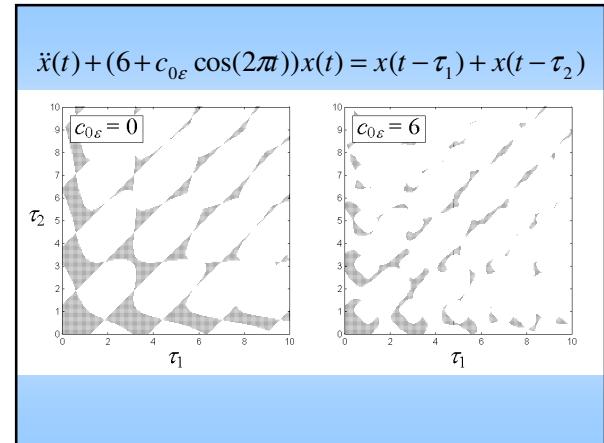
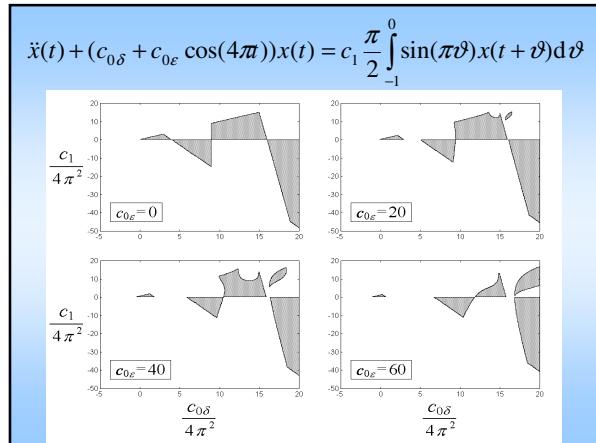
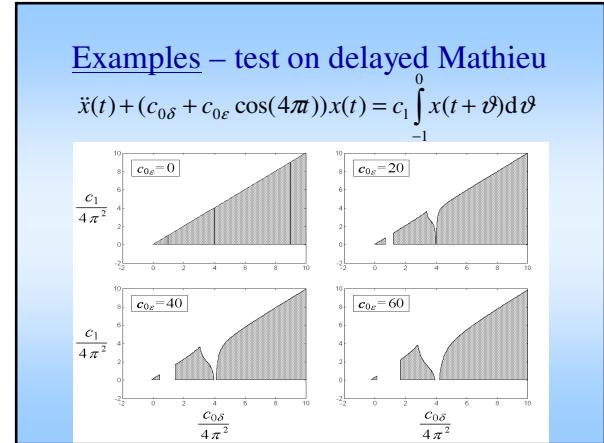
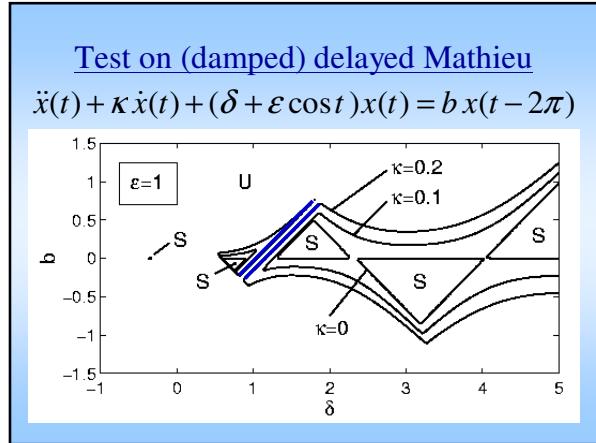
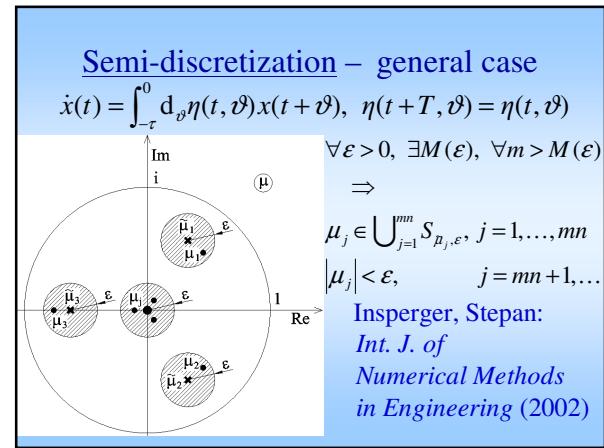
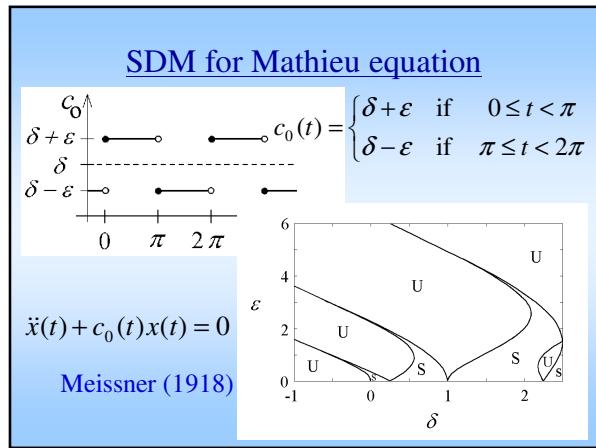
$$\mathbf{y}_i = \begin{pmatrix} x(t_i) \\ \dot{x}(t_i) \end{pmatrix} \quad \mathbf{A}_i = \begin{pmatrix} \cos(\Delta t \sqrt{c_{0i}}) & \frac{\sin(\Delta t \sqrt{c_{0i}})}{\sqrt{c_{0i}}} \\ -\sqrt{c_{0i}} \sin(\Delta t \sqrt{c_{0i}}) & \cos(\Delta t \sqrt{c_{0i}}) \end{pmatrix}$$

$$\begin{pmatrix} x(2\pi) \\ \dot{x}(2\pi) \end{pmatrix} = \mathbf{y}_k = \mathbf{A}_{k-1} \mathbf{A}_{k-2} \dots \mathbf{A}_0 \mathbf{y}_0 = \Phi(2\pi) \begin{pmatrix} x(0) \\ \dot{x}(0) \end{pmatrix}$$

$$\det(\mu \mathbf{I} - \Phi(2\pi)) = 0 \Rightarrow |\mu_{1,2}| \leq 1 \Leftrightarrow \text{stability}$$

for $k=2$ intervals \Rightarrow Meissner (1918)

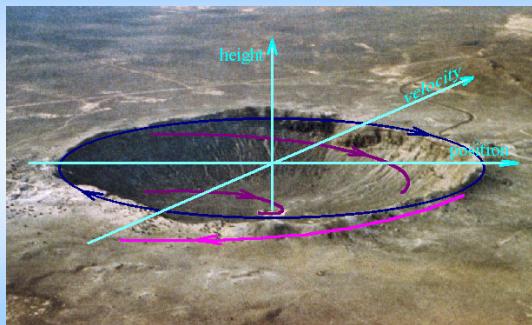
for $k \rightarrow \infty \Rightarrow$ van der Pol, Strutt, Ince (1928)



Nonlinear RFDEs in Engineering

Stability analysis of steady-states is followed by
 Center Manifold reduction & bifurcation analysis
 Hopf bifurcation – self-excited vibrations
Supercritical case: easy to avoid vibrations by
 knowing the linear stability behavior
Subcritical case: the unstable periodic solutions
 mean a limited domain of attraction for the
 desired steady-state behavior – *cannot* be
 predicted by linear stability analysis.

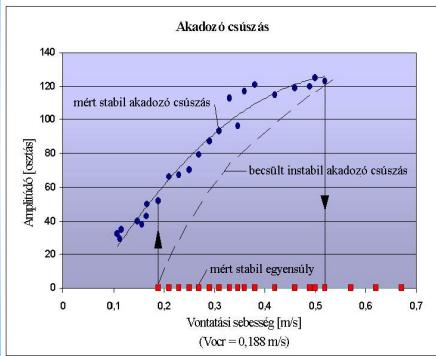
Unstable limit cycle – “ghost” vibration



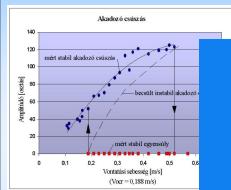
Stick&slip



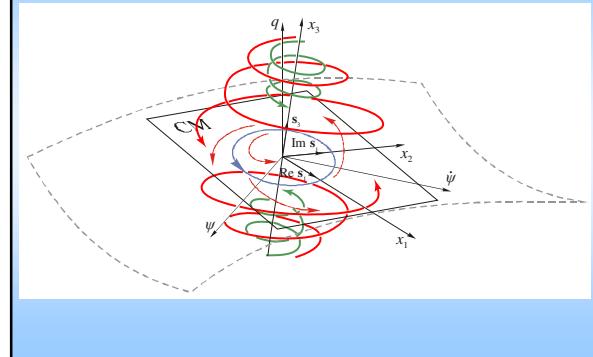
Stick&slip experimental bifurcation diagram



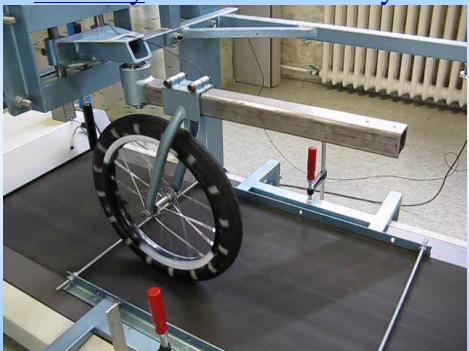
Stick&slip experimental bifurcation diagram



Center Manifold reduction



Shimmy – unstable limit cycle



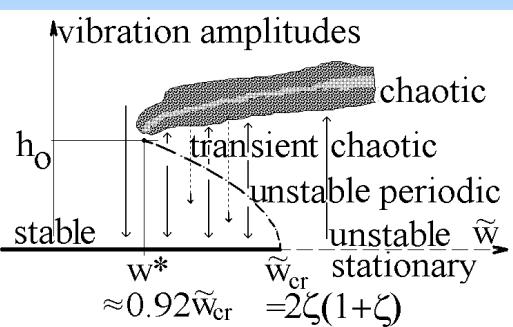
Shimmy – quasi-periodic oscillations



Machine tool vibrations

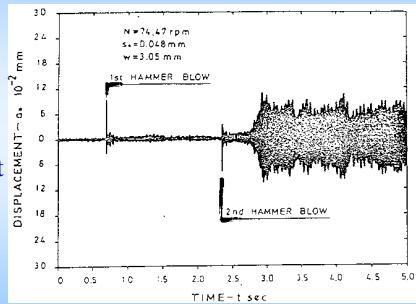


Machine tool vibration



The unstable periodic motion

Shi, Tobias
(1984) –
impact
experiment



Balancing

